

LA-UR-20-26424

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Title: Searching for new physics with single and double beta decay

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Intended for: LANL P/T Colloquium

Issued: 2020-08-20

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Searching for new physics with single and double beta decay

Vincenzo Cirigliano

Los Alamos National Laboratory



U.S. DEPARTMENT OF
ENERGY

Office of
Science



Outline

- The quest for new physics: Energy and Precision Frontiers
- Two Precision Frontier probes with strong LANL involvement:
 - Precision β decay measurements
 - Neutrinoless $\beta\beta$ decay

Acknowledgements:

T-2 / CCS-7 collaborators: T. Bhattacharya, K. Fuyuto, M. Graesser, R. Gupta, E. Mereghetti, B. Yoon

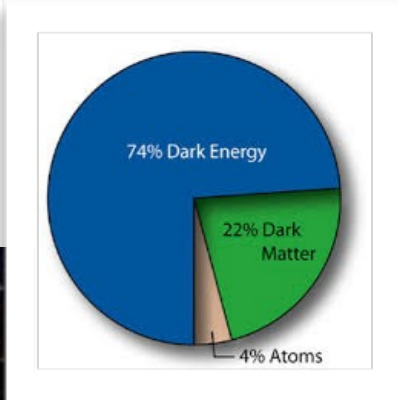
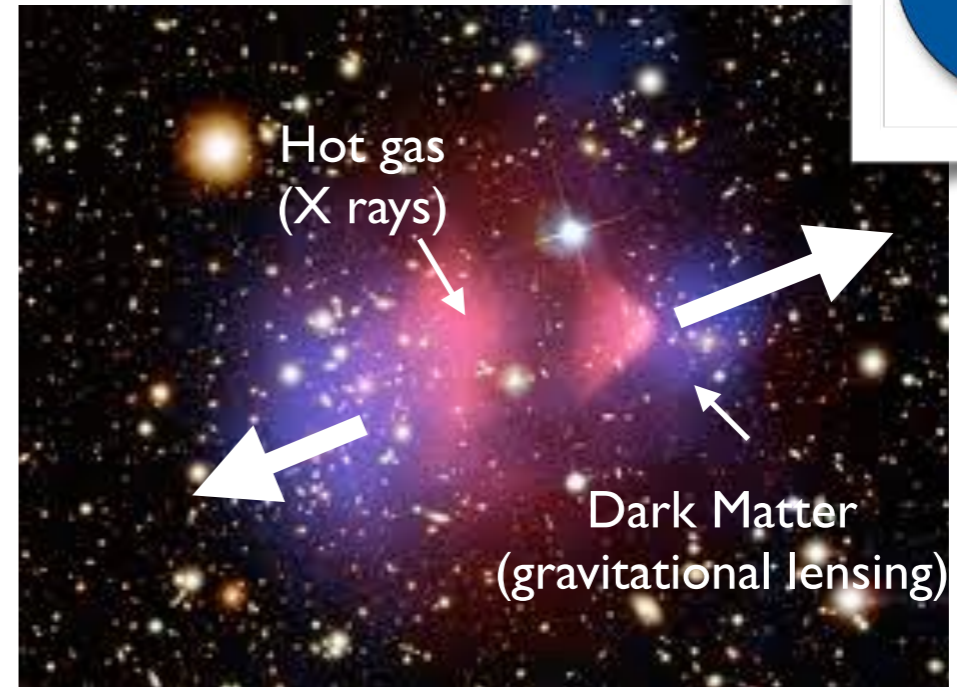
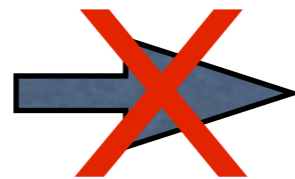
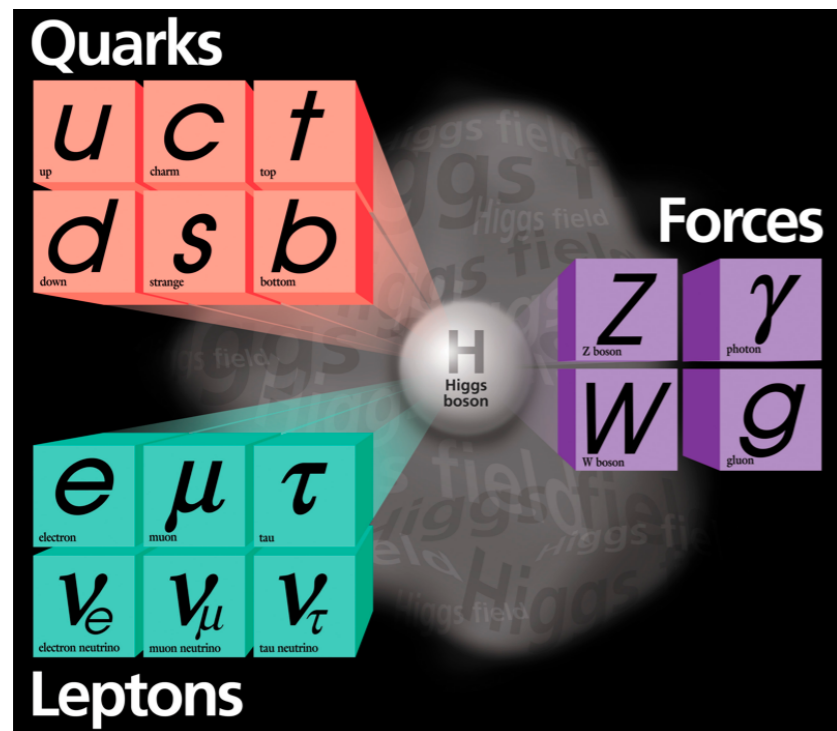
External collaborators: W. Dekens, J. de Vries, S. Pastore, M. Piarulli, R. Wiringa, U. Van Kolck

P-23 Weak Interactions Team

P-25 Neutron Physics Team

The quest for new physics

New physics: why?



No Baryonic Matter, no Dark Matter, no Dark Energy, no Neutrino Mass

What stabilizes $G_{\text{Fermi}}/G_{\text{Newton}}$ against radiative corrections?

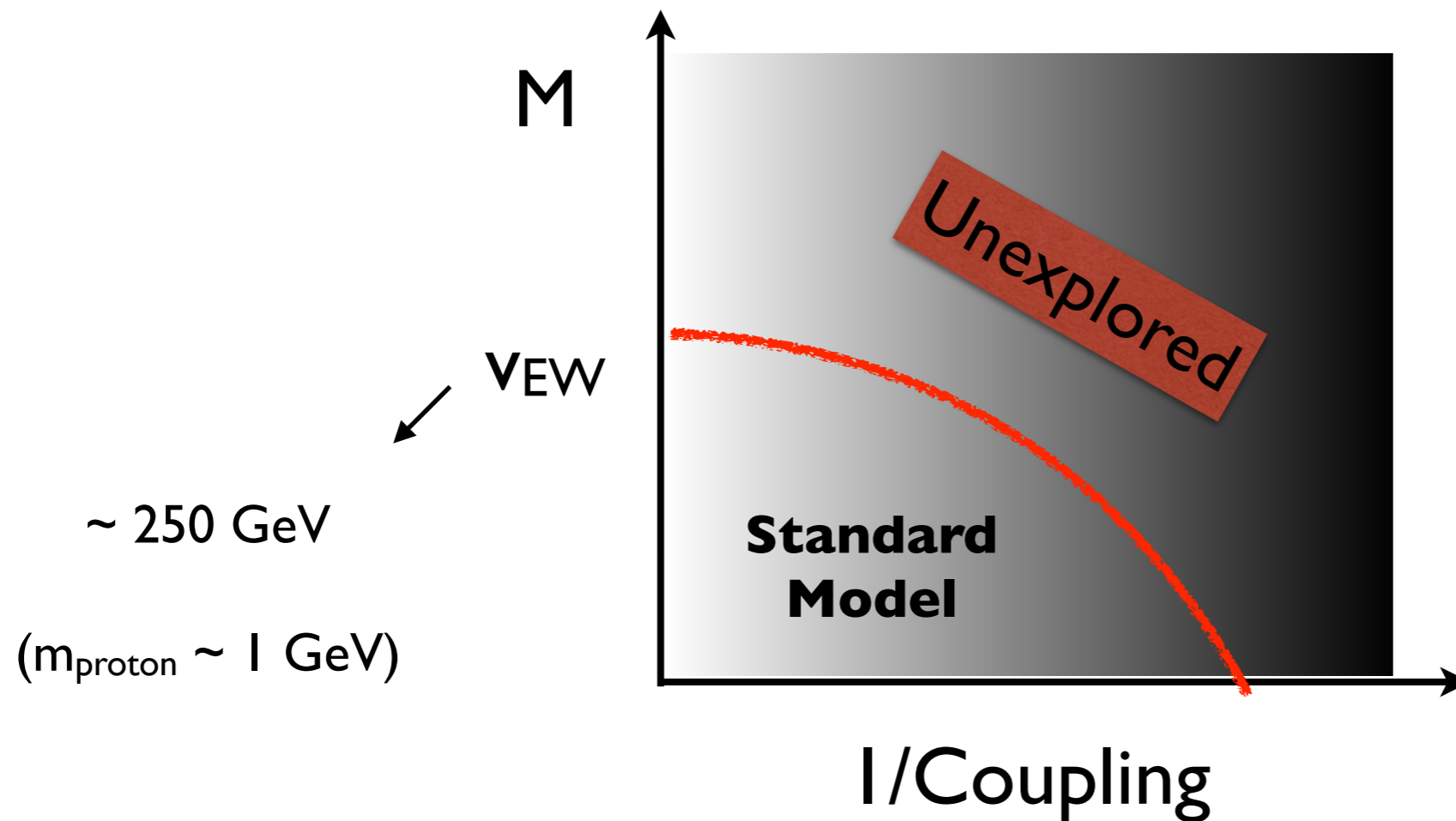
Do forces unify at high E? What is the origin of families?

...

Addressing these puzzles likely requires new degrees of freedom

New physics: where?

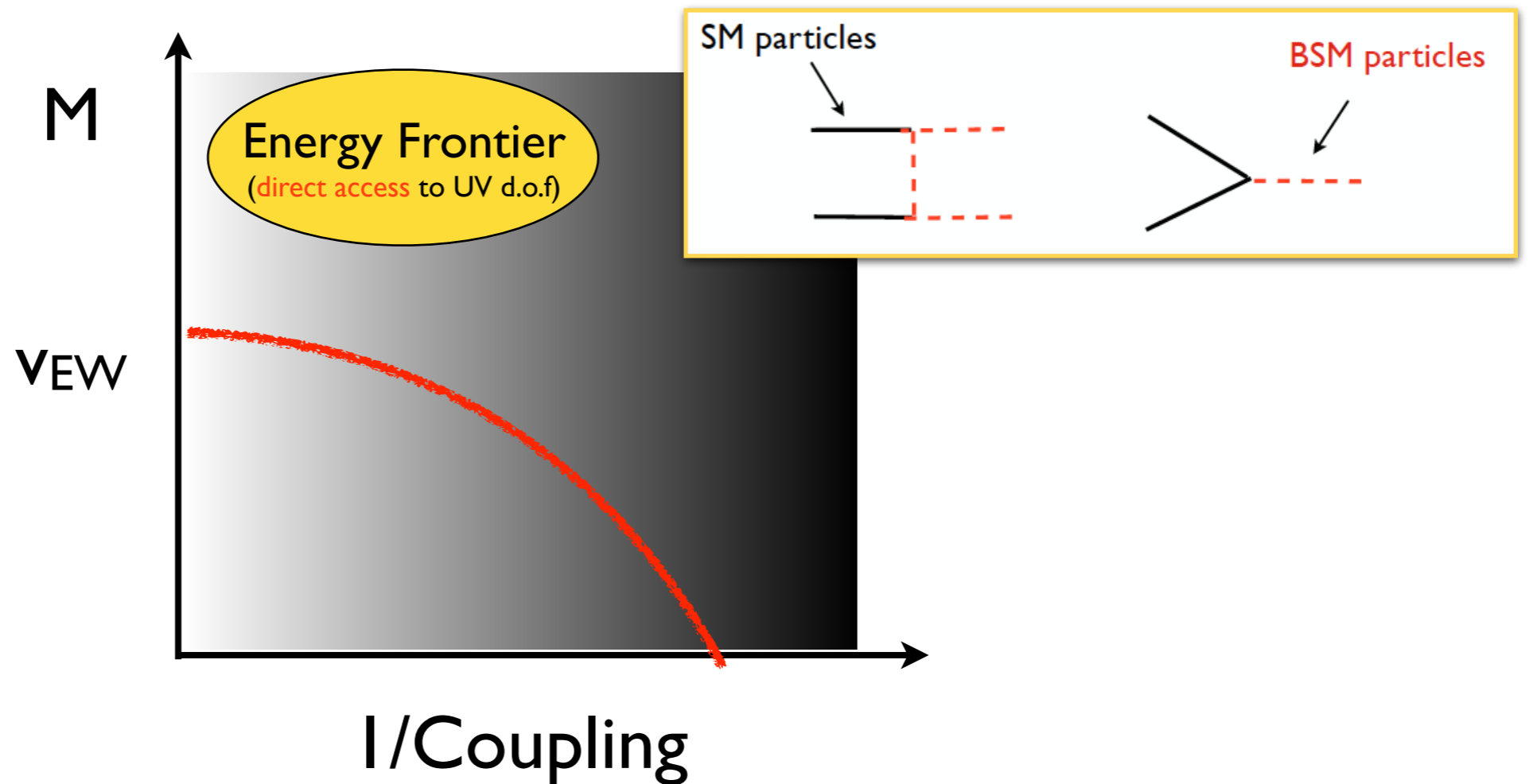
- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?



New physics: how?

- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?

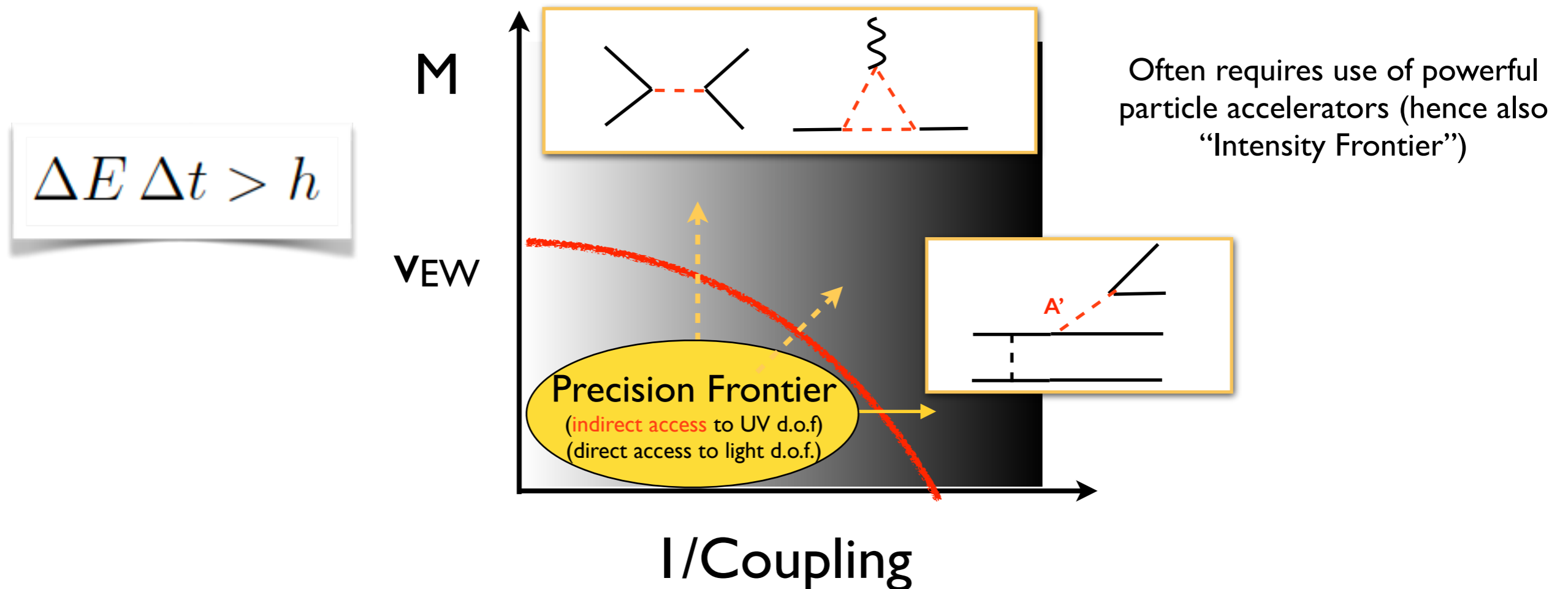
$$E = mc^2$$



- Two approaches

New physics: how?

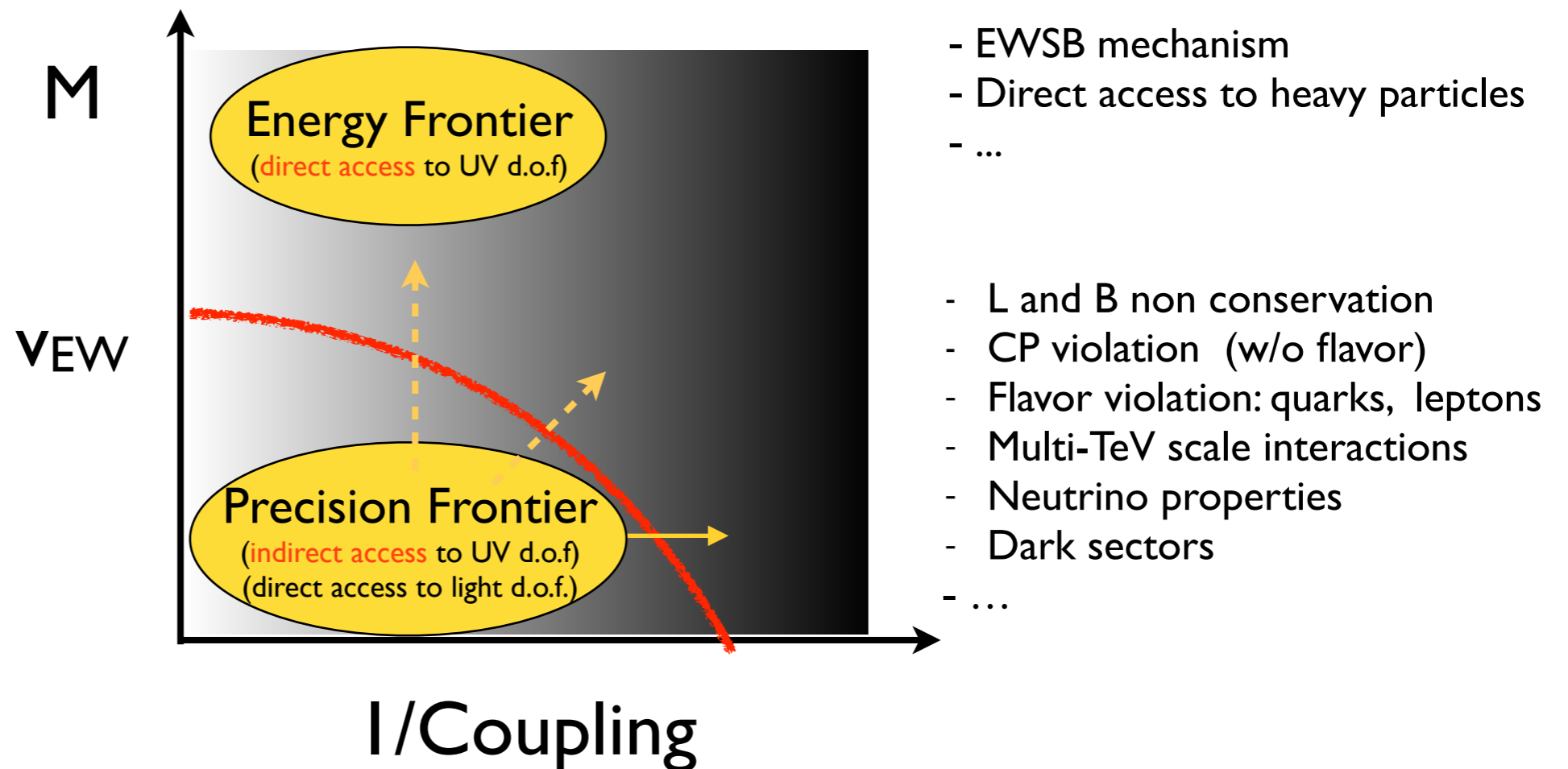
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- Two approaches

New physics: how?

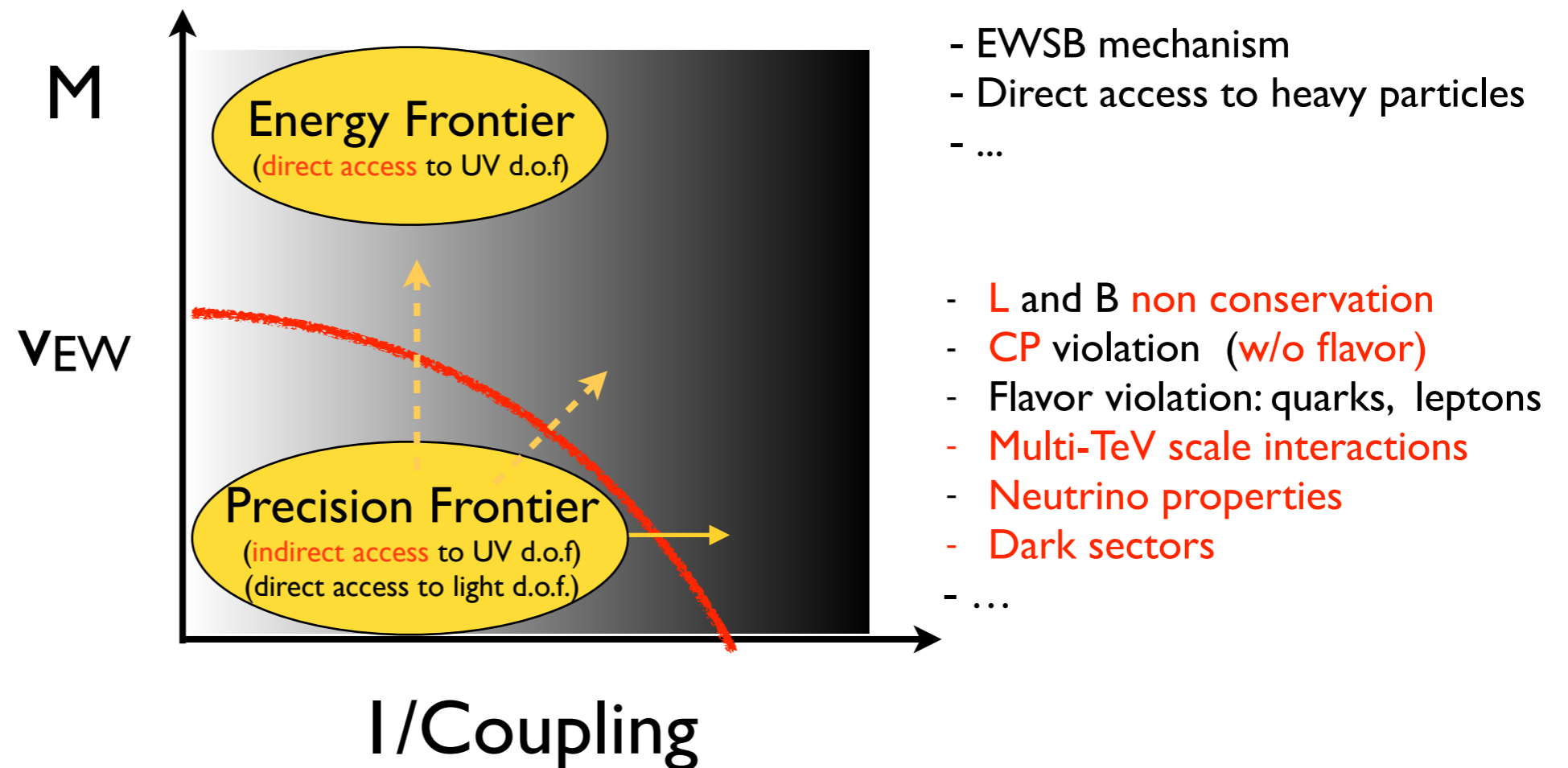
- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?



- Two approaches, both needed to reconstruct BSM dynamics:
structure, symmetries, and parameters of \mathcal{L}_{BSM}

New physics: how?

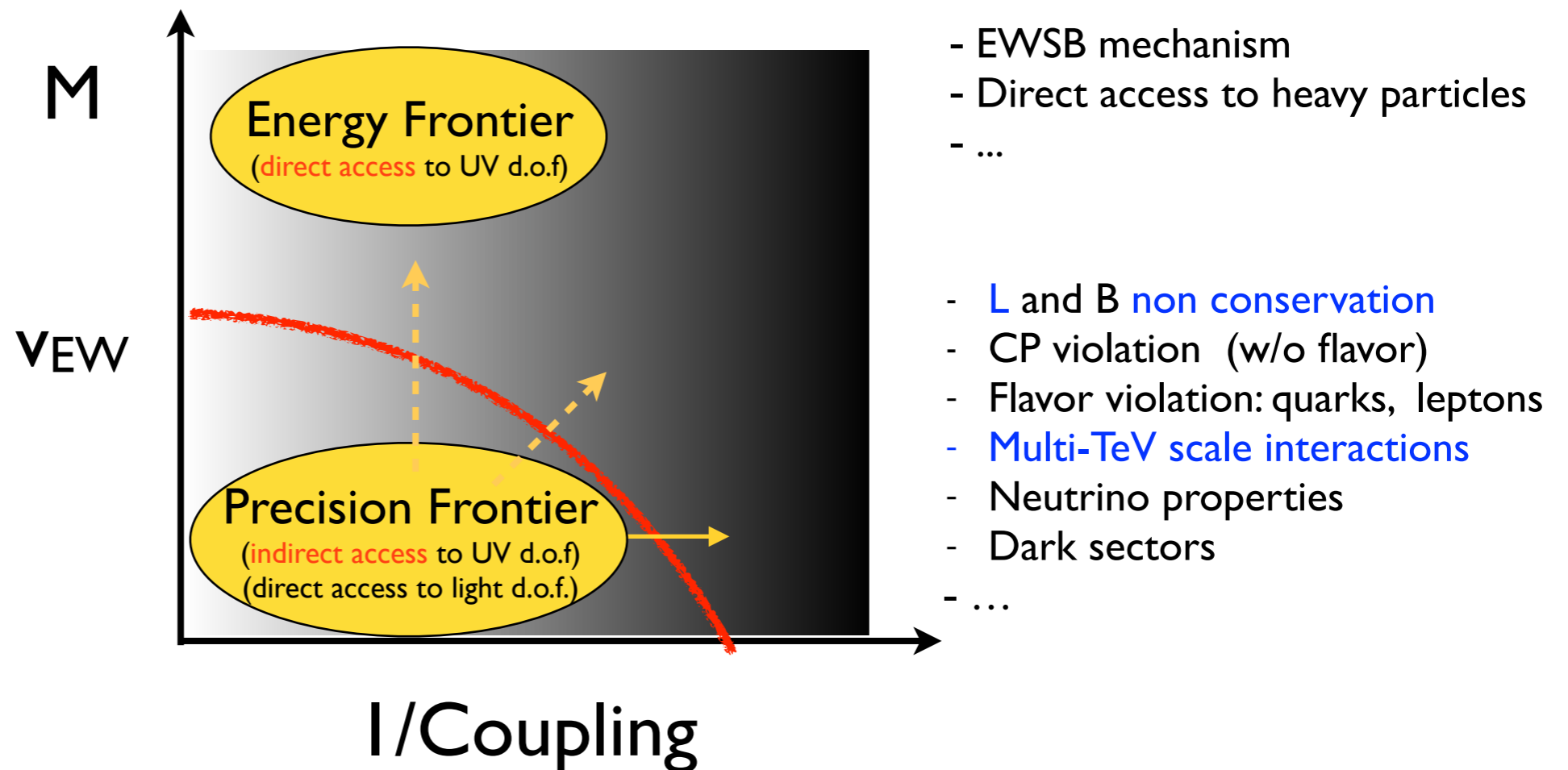
- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?



LANL NP & HEP programs
play a prominent role at the Precision Frontier

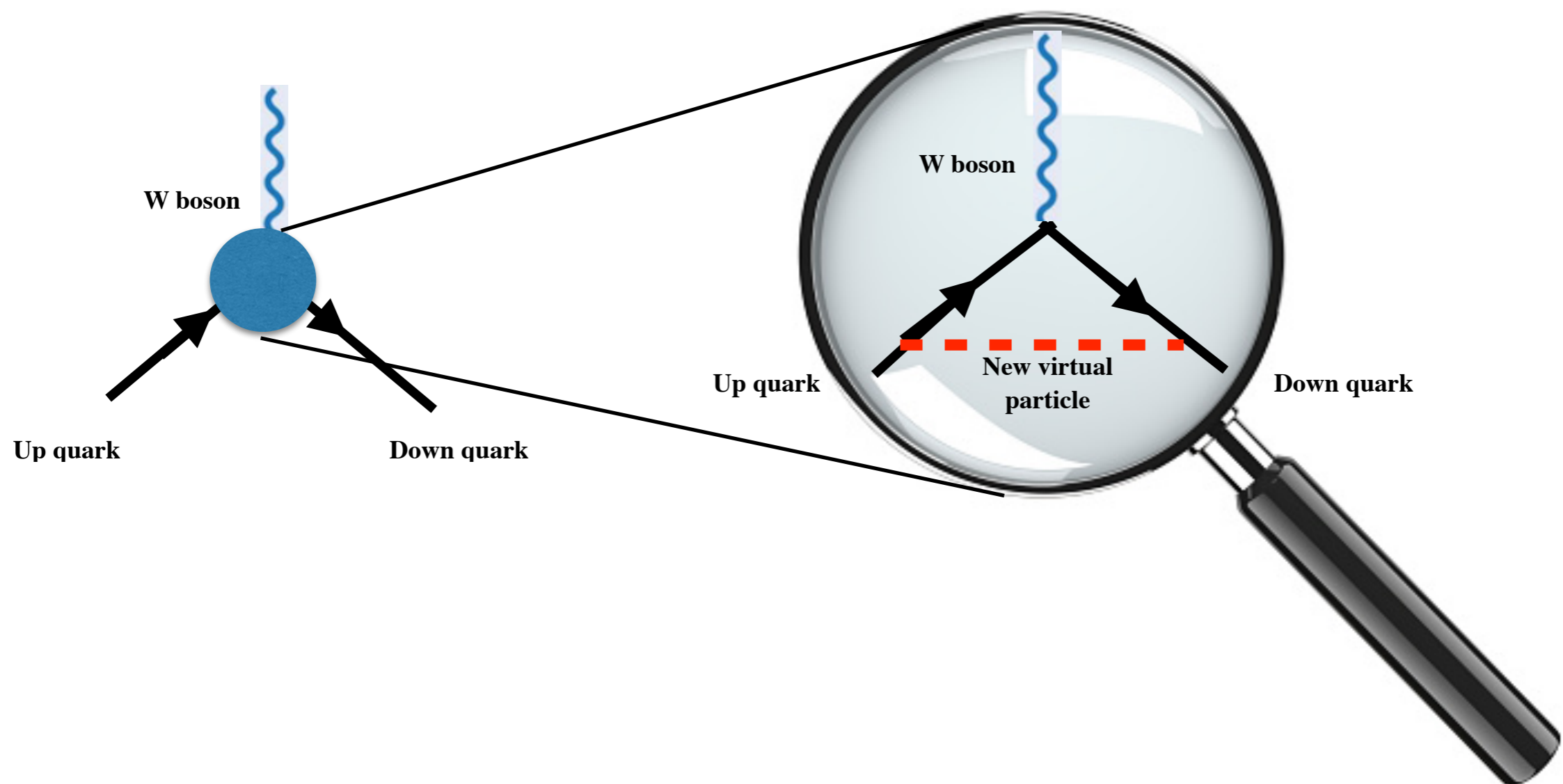
New physics: how?

- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?

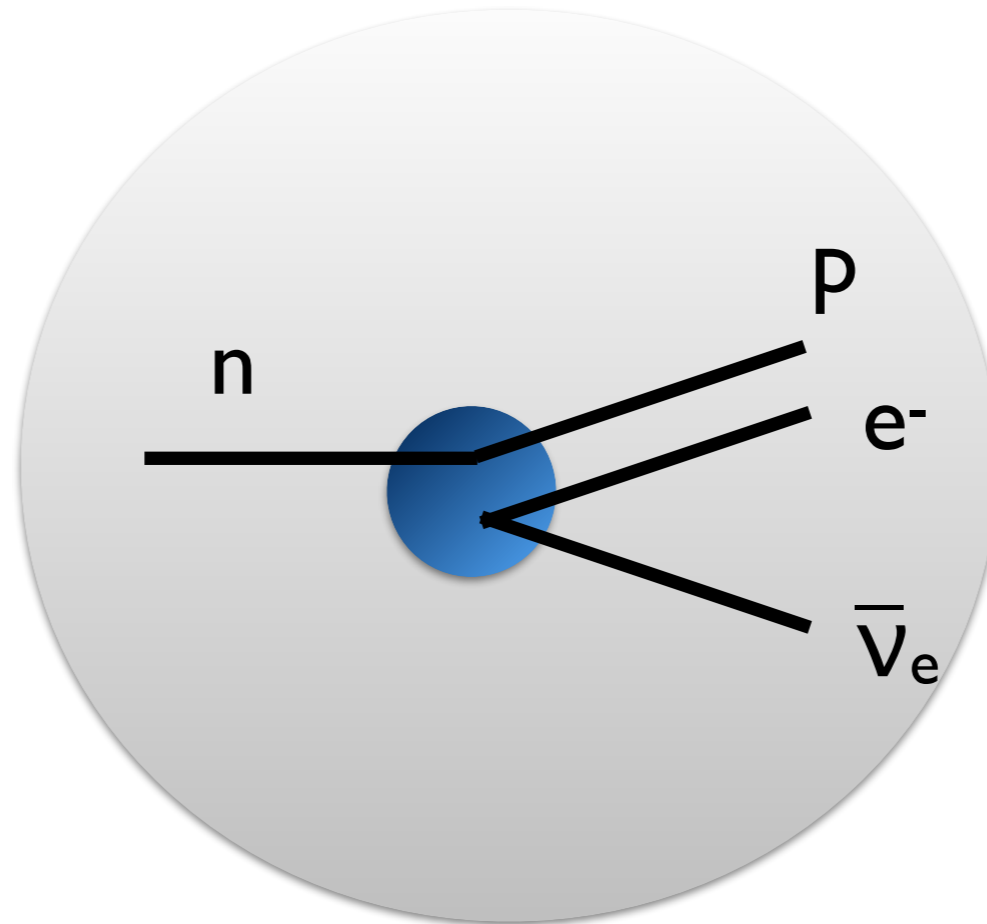


I will discuss β and $\beta\beta$ decays as probes of BSM weak interactions and $L\#$ non-conservation, respectively

Beta decays as a probe of new physics



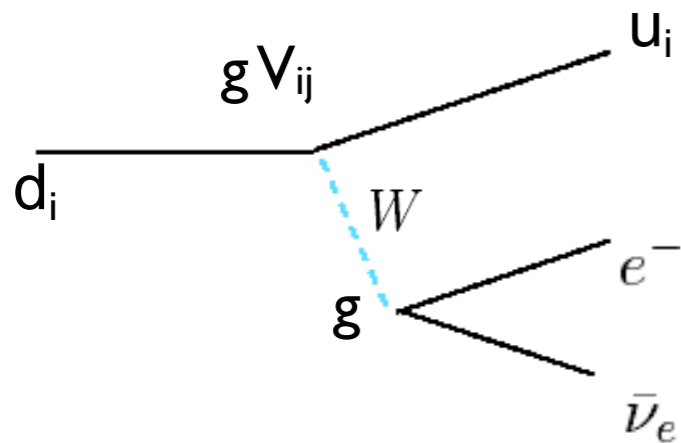
Beta decays in the SM and beyond



- Beta decays have played a central role in the development of the SM
- Nowadays: tool to challenge the SM & probe possible new physics

Beta decays in the SM and beyond

- In the SM, W exchange \Rightarrow V-A currents, universality



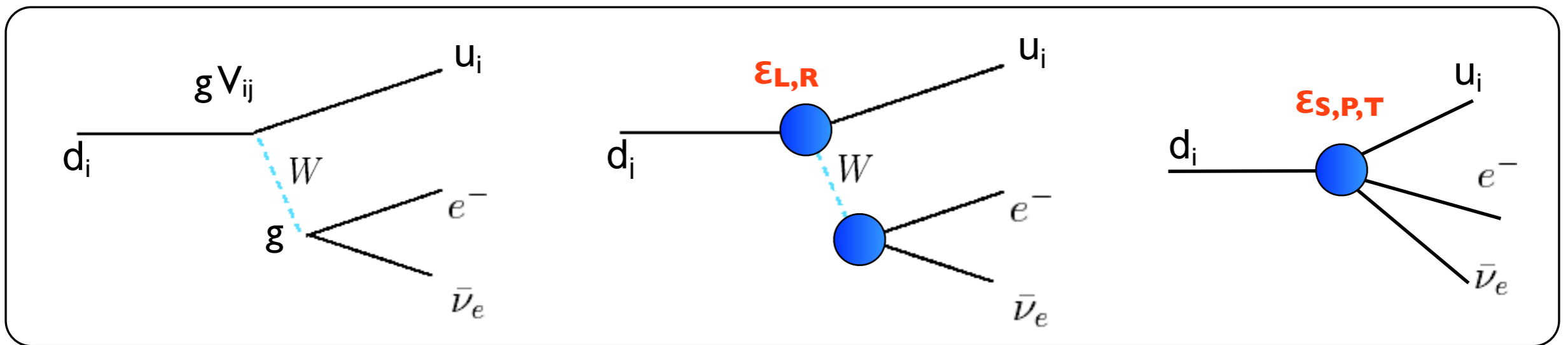
$$G_F^{(\beta)} \sim g^2 V_{ij} / M_W^2 \sim G_F^{(\mu)} V_{ij} \sim 1/v^2 V_{ij}$$

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Cabibbo-Kobayashi-Maskawa

Beta decays in the SM and beyond

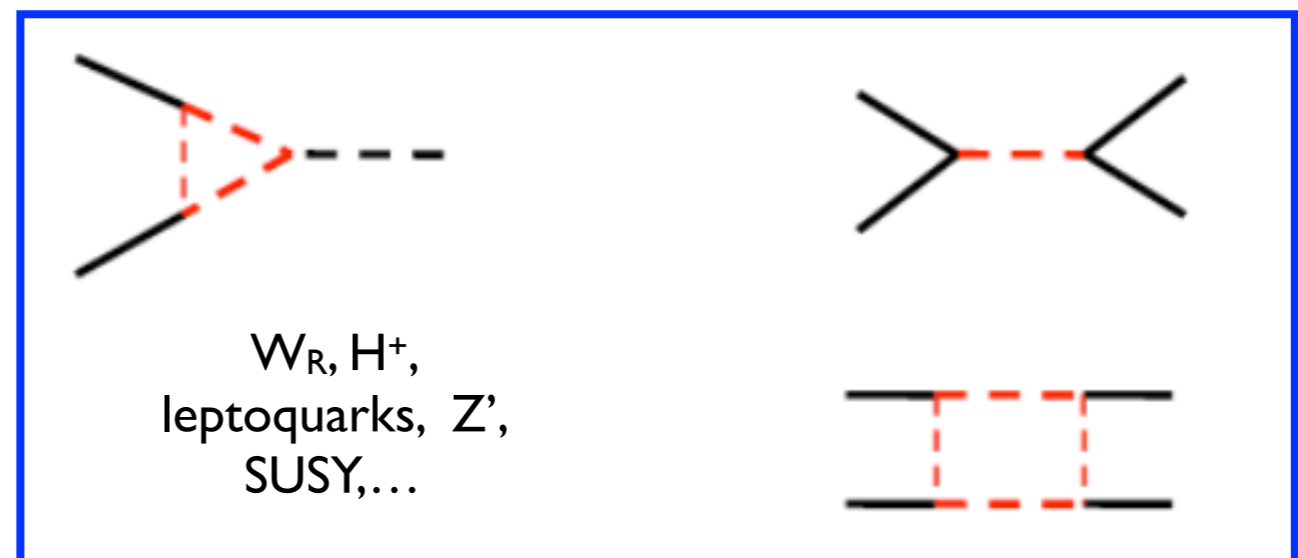
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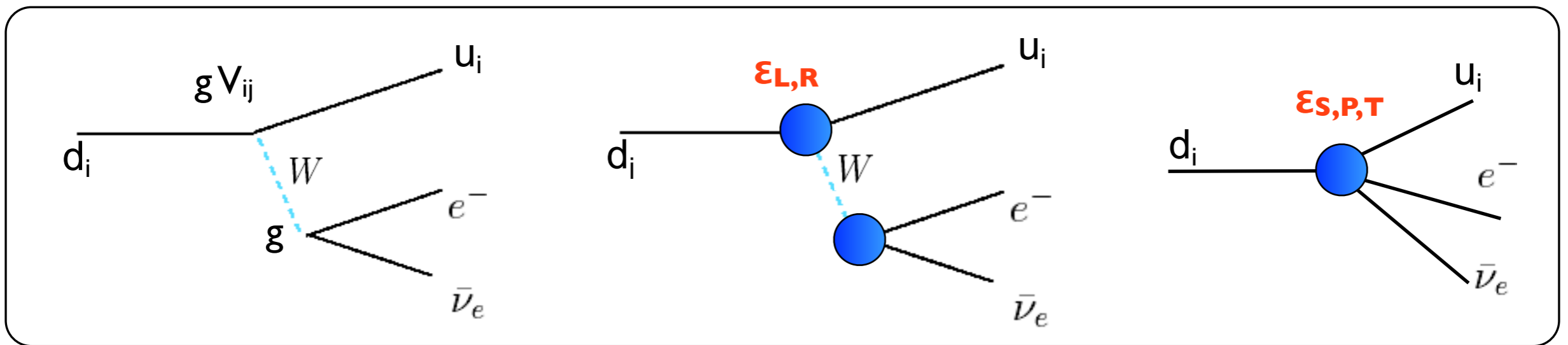
$$1/\Lambda^2$$

$$1/\Lambda^2$$



Beta decays in the SM and beyond

- In the SM, W exchange \Rightarrow V-A currents, universality



$$G_F(\beta) \sim g^2 V_{ij} / M_W^2 \sim G_F(\mu) V_{ij} \sim 1/v^2 V_{ij}$$

$$1/\Lambda^2$$

$$1/\Lambda^2$$

$$E \ll \Lambda \quad \downarrow \quad \epsilon_\Gamma \sim \tilde{\epsilon}_\Gamma \sim (v/\Lambda)^2$$

$$\mathcal{L}_{\text{SM}} = \frac{G_F V_{ud}}{\sqrt{2}} \sum_{\Gamma} \left[\epsilon_\Gamma \bar{\ell} \Gamma \nu_L \cdot \bar{u} \Gamma d + \tilde{\epsilon}_\Gamma \bar{\ell} \Gamma \nu_R \cdot \bar{u} \Gamma d \right]$$

Ten effective couplings

$$\Gamma = L, R, S, P, T$$

How do we probe the ε_α ? (I)

I. Differential decay distribution (mostly sensitive to $\varepsilon_{S,T}$)

$$d\Gamma \propto F(E_e) \left\{ 1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \langle \vec{J} \rangle \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + \dots \right] \right\}$$

Lee-Yang, 1956 Jackson-Treiman-Wyld 1957

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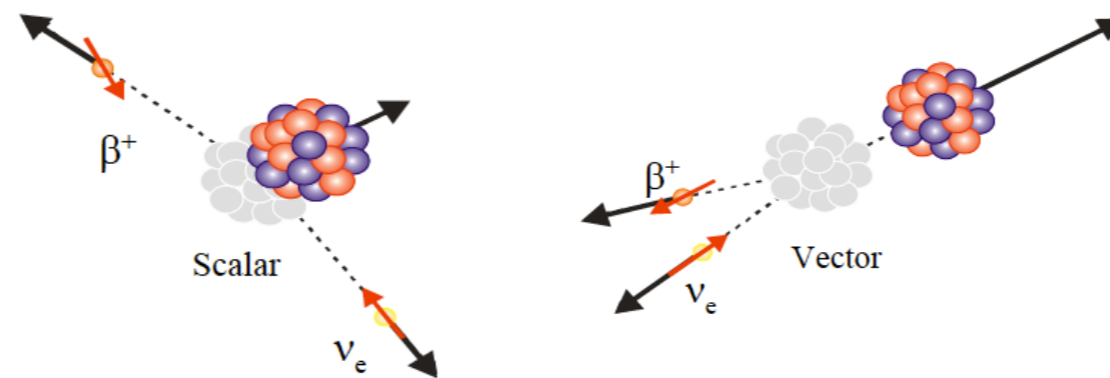
b ($g_S \varepsilon_S$, $g_T \varepsilon_T$):
distortion of beta spectrum

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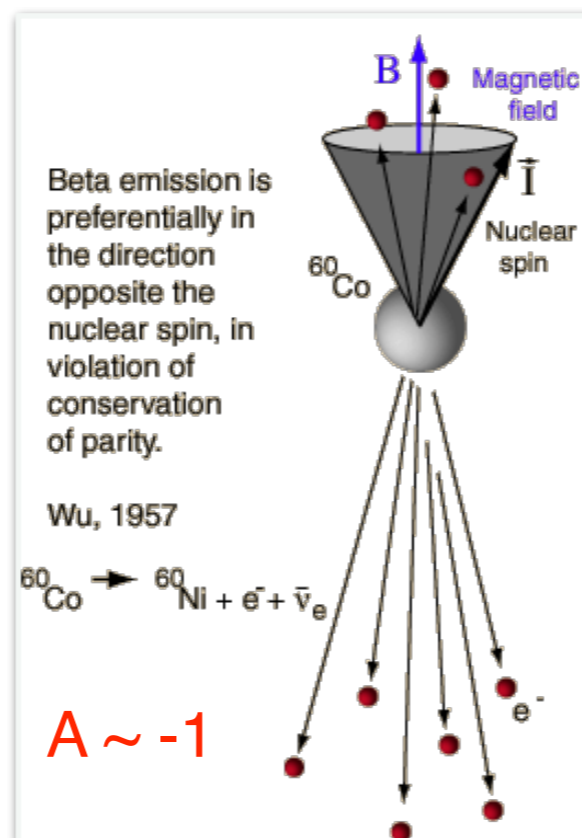
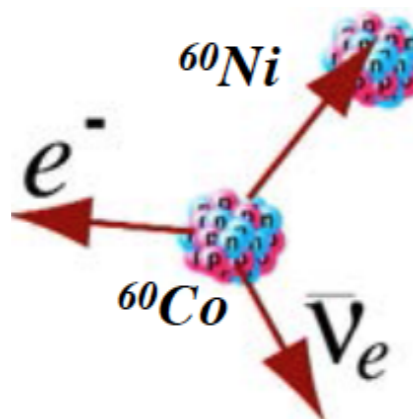
How do we probe the ϵ_α ? (I)

I. Differential decay distribution (mostly sensitive to $\epsilon_{S,T}$)

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Lee-Yang, 1956

Jackson-Treiman-Wyld 1957



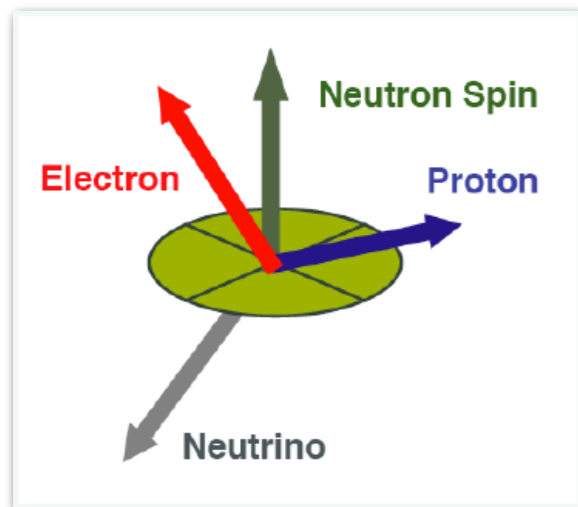
C-S Wu

How do we probe the ε_α ? (I)

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$$d\Gamma \propto F(E_e) \left\{ 1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \langle \vec{J} \rangle \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + \dots \right] \right\}$$

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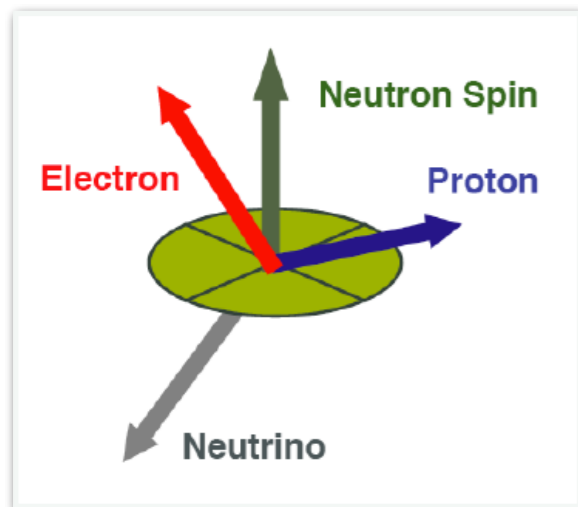
$a(g_A)$, $A(g_A)$, $B(g_A, g_\alpha \varepsilon_\alpha)$, ...
isolated via suitable experimental
asymmetries

How do we probe the ε_α ? (I)

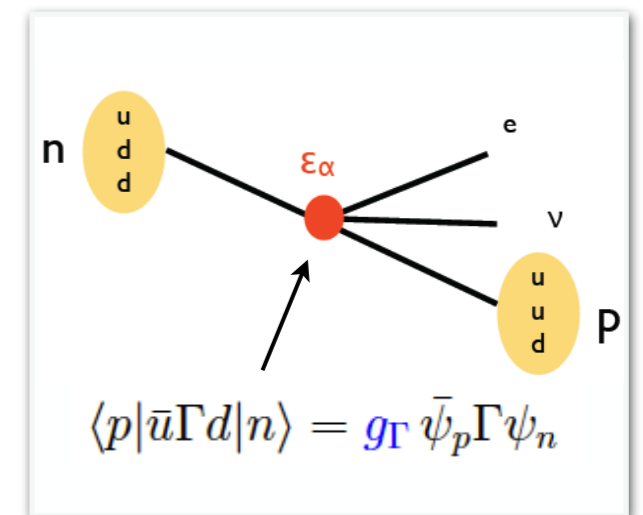
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$a(g_A)$, $A(g_A)$, $B(g_A, g_\alpha \varepsilon_\alpha)$, ...
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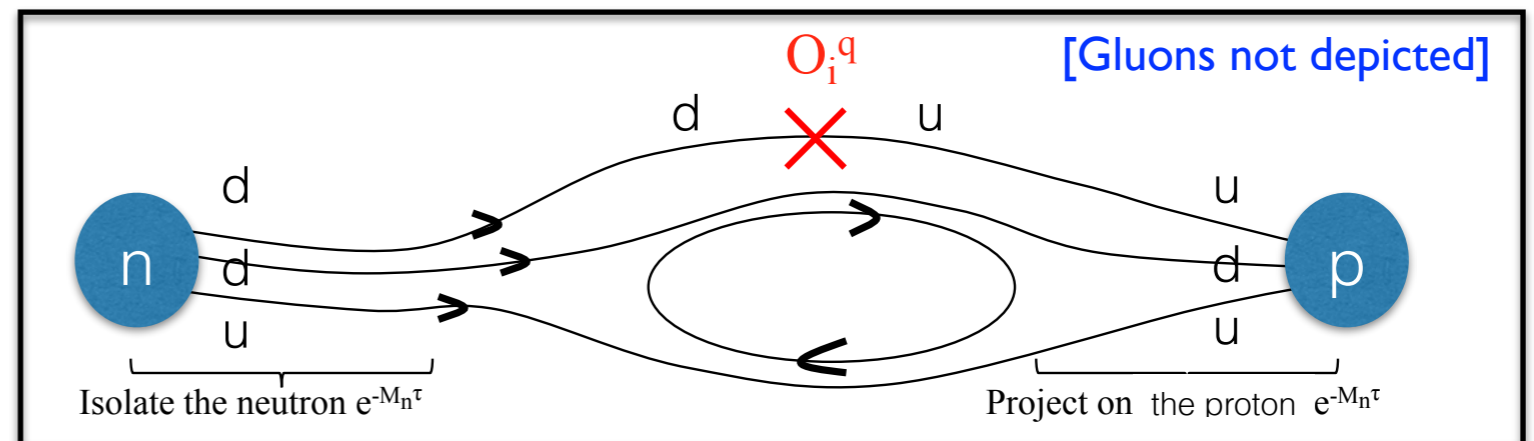
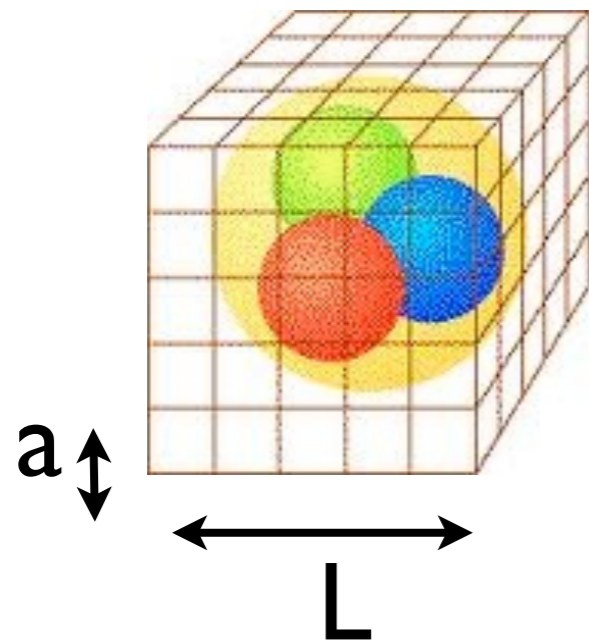


Theory input: nucleon charges $g_{A,S,T}$

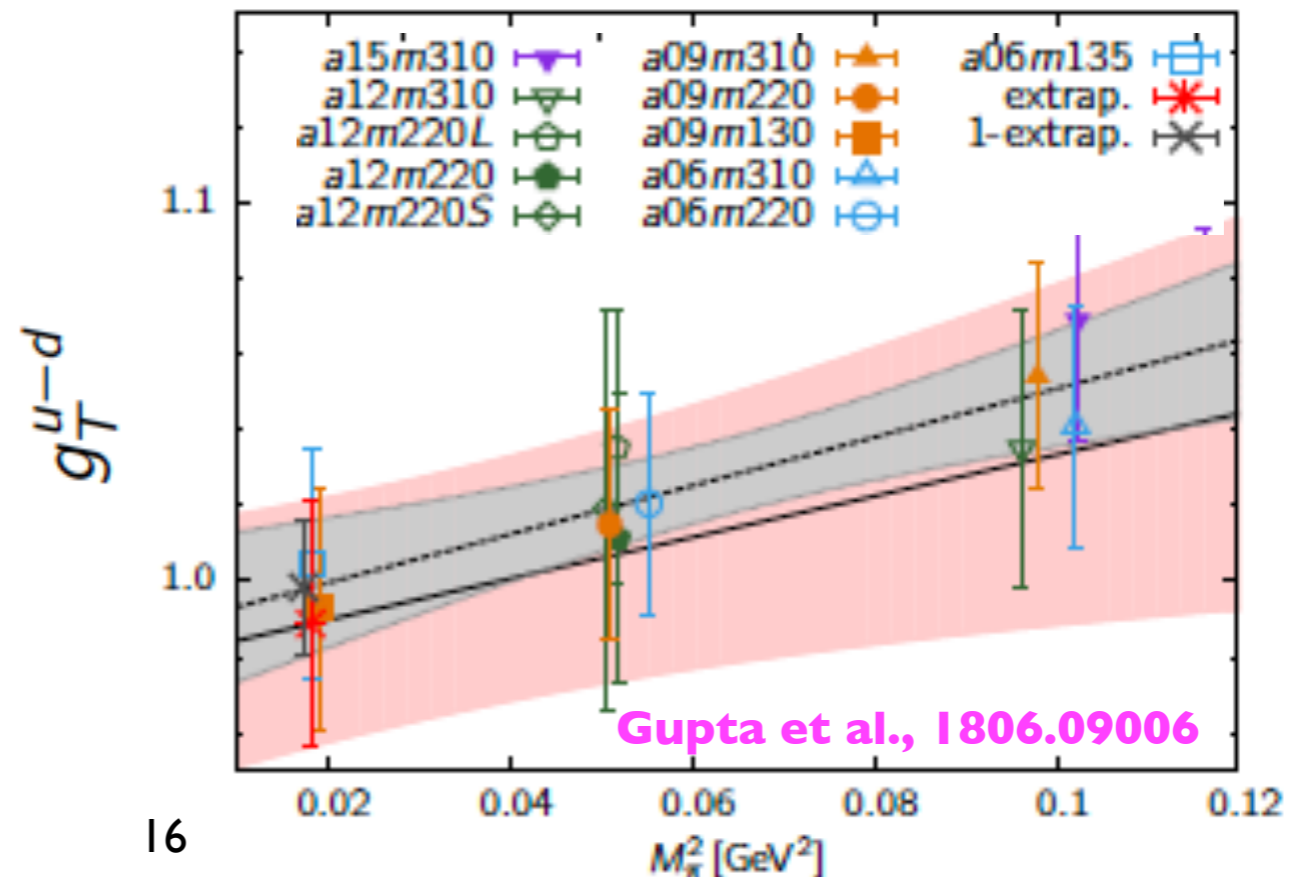
Great progress in lattice QCD, spearheaded by LANL theory group

Nucleon charges from lattice QCD

- Discretize space-time into a finite Euclidean lattice $(a, V) \rightarrow$ perform Monte Carlo integration of the path integral

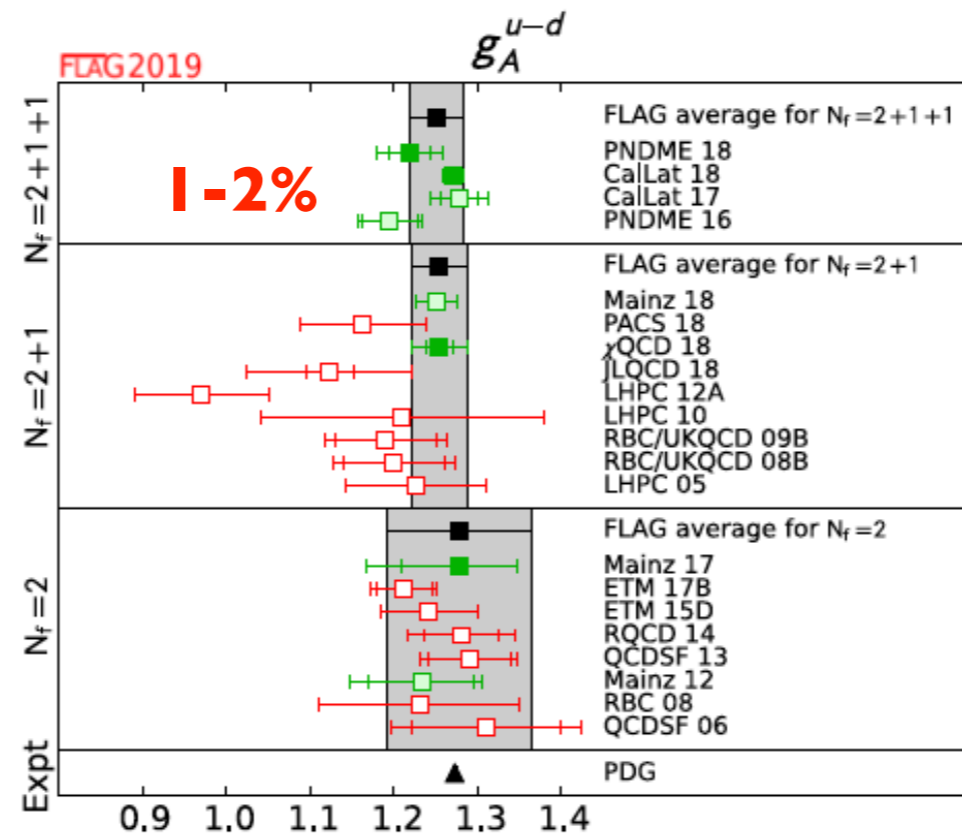


- Remove all systematics by performing calculation on many little universes with different m_q , a , L

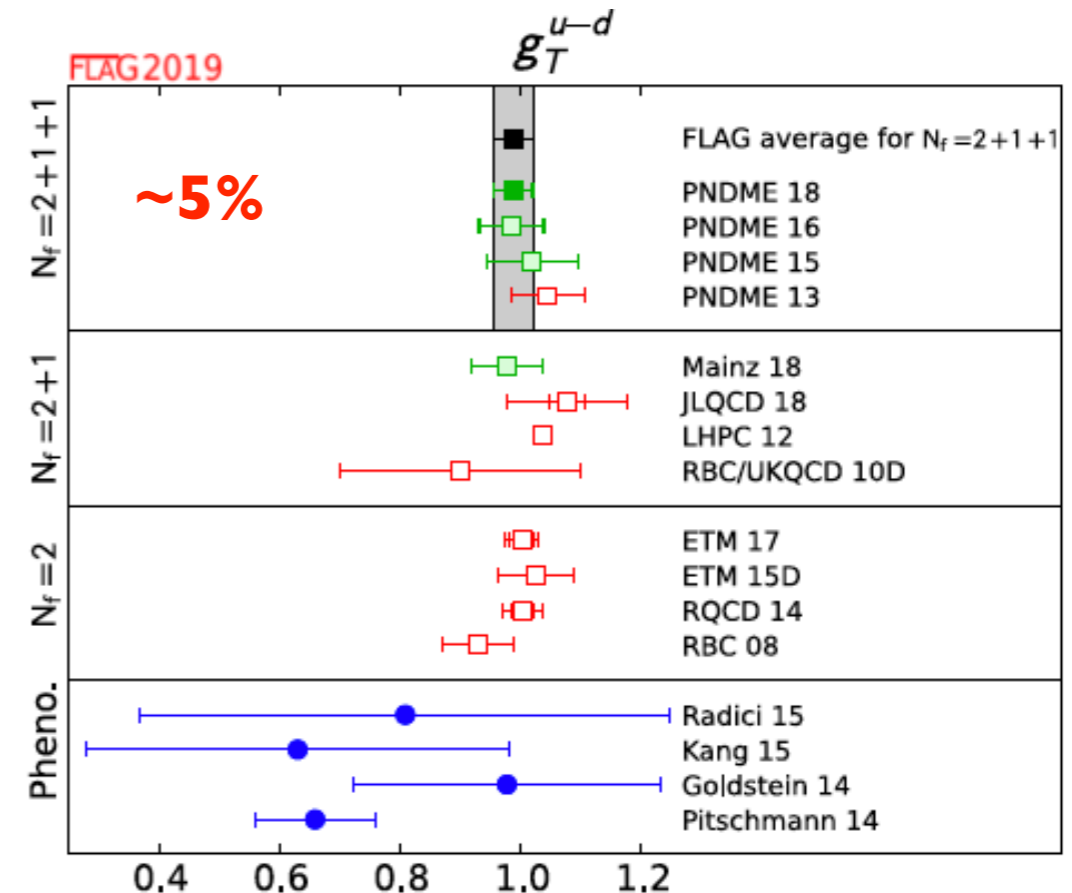
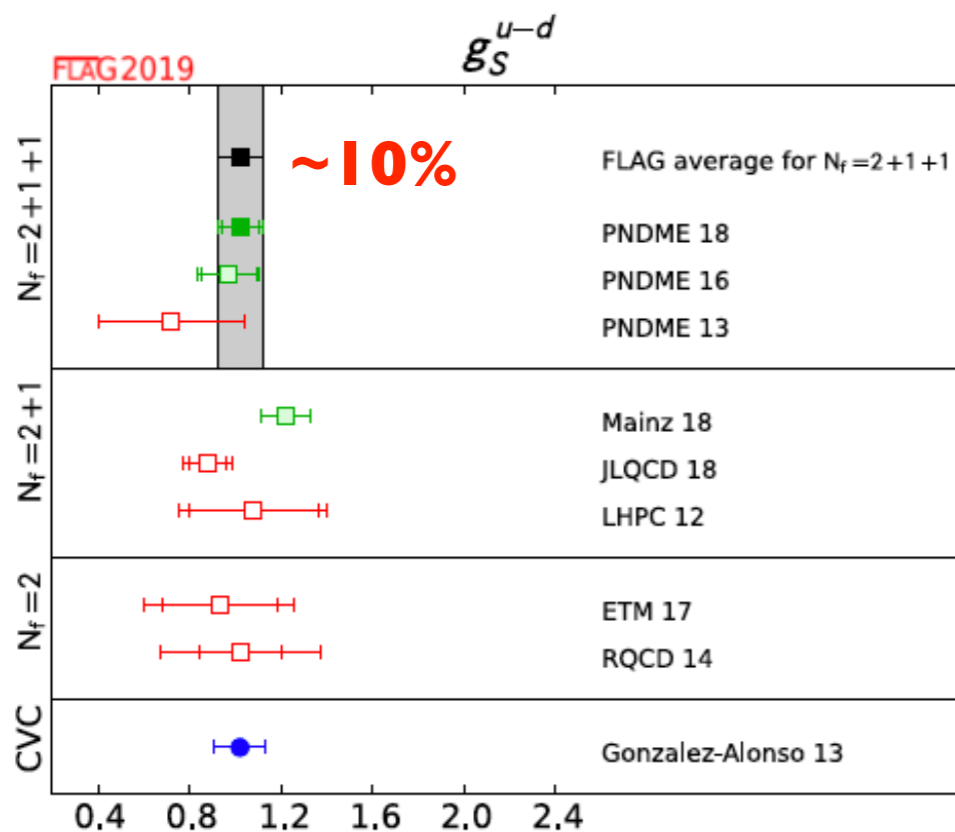


Nucleon charges from lattice QCD

World averages dominated by LANL results

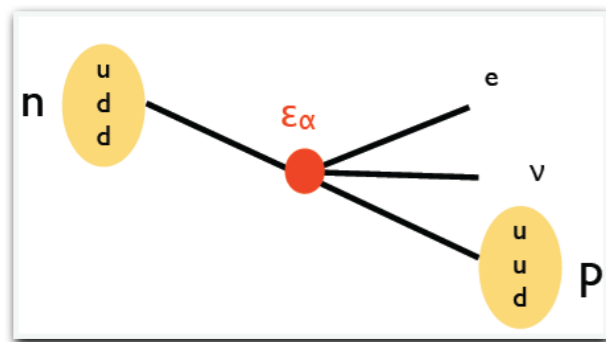


FLAG review
1902.08191



Sensitivity to ϵ_S and ϵ_T

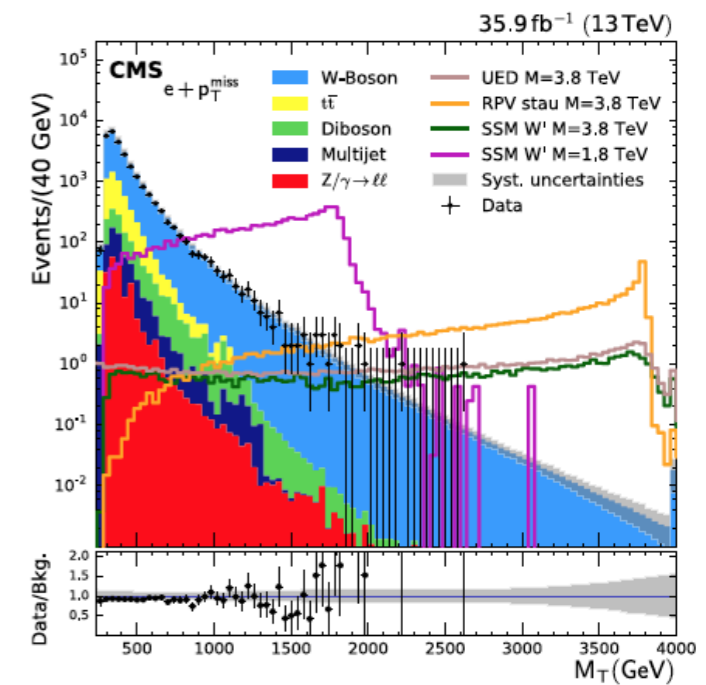
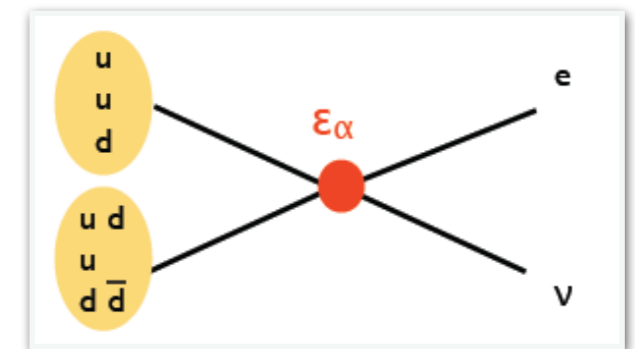
$$n \rightarrow p e \nu$$



Current low-E data:
dominated by
 $b(0^+ \rightarrow 0^+)$, $A(n)$

**Gonzalez-Alonso,
Naviliat-Cuncic,
Severijns, 1803.08732**

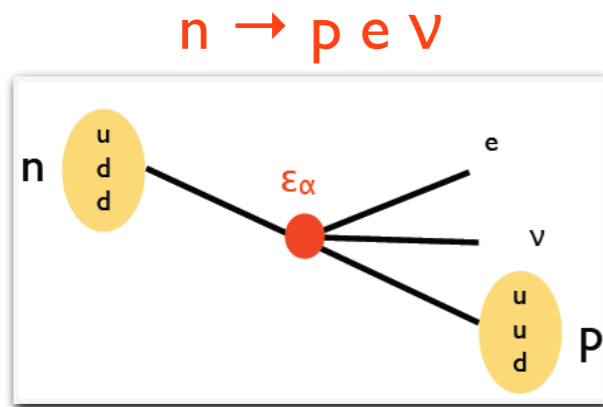
$$\text{LHC: } pp \rightarrow e \nu + X$$



Sensitivity to ϵ_S and ϵ_T

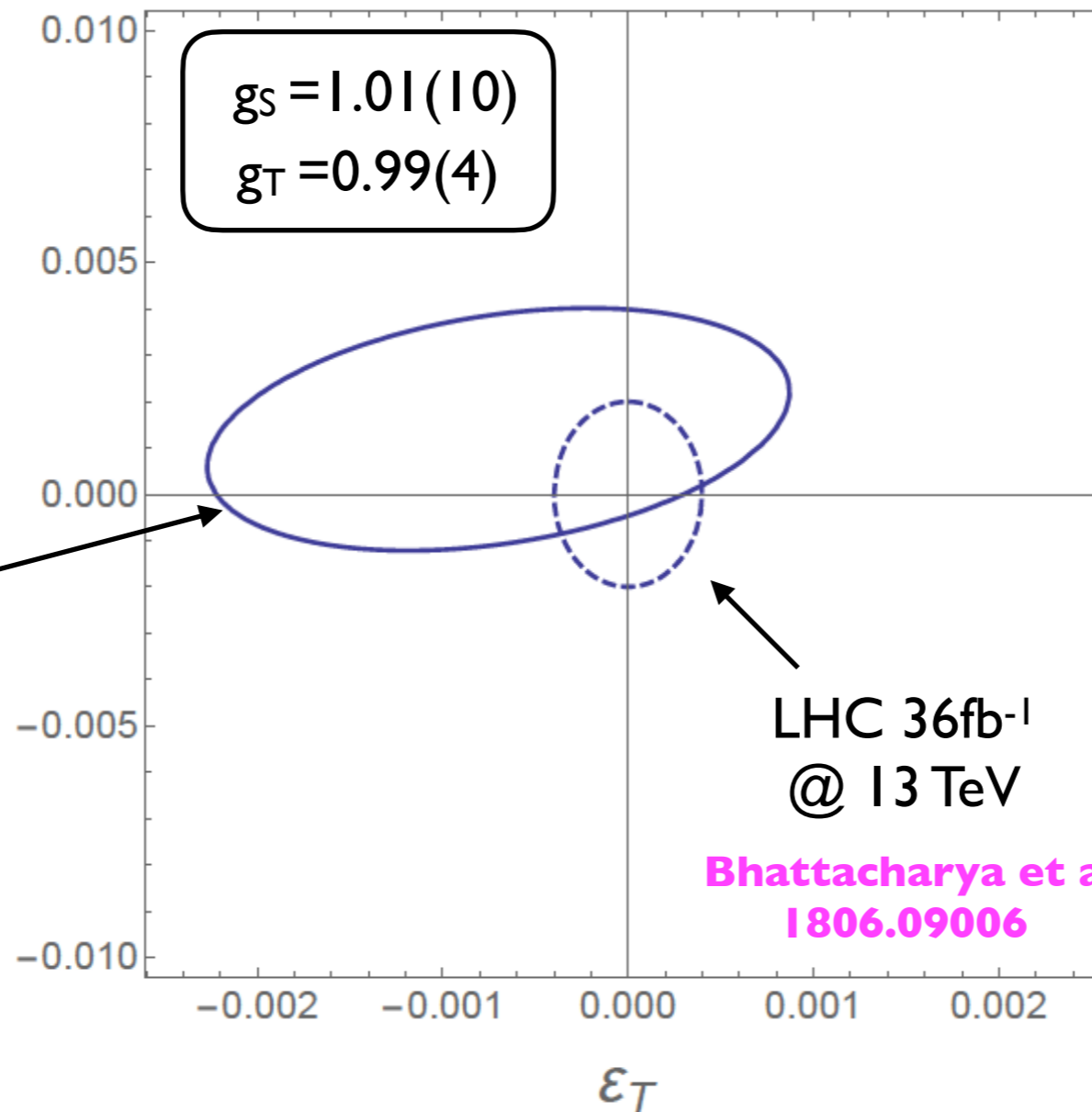
CURRENT

$\epsilon_{S,T}$ @ $\mu = 2 \text{ GeV (MS-bar)}$

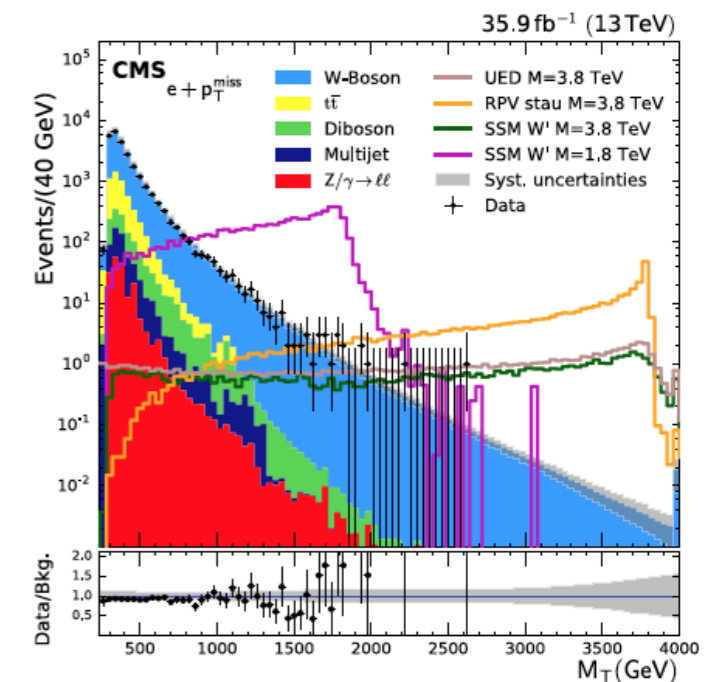
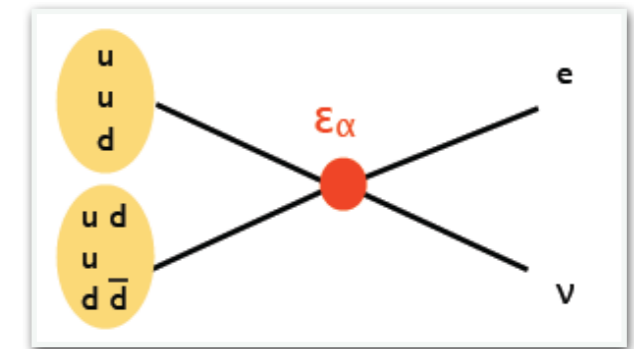


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Gonzalez-Alonso,
Naviliat-Cuncic,
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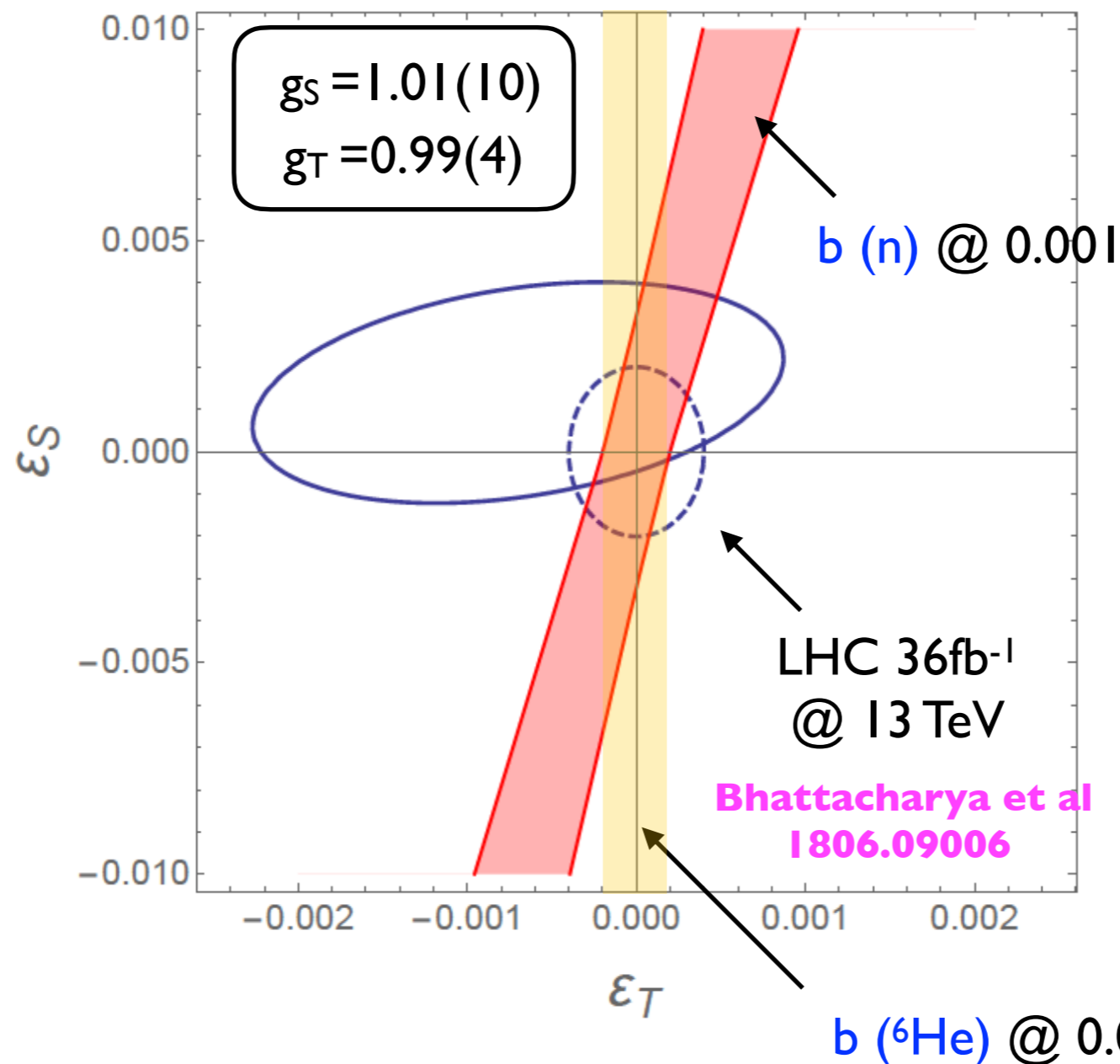
LHC: $pp \rightarrow e \nu + X$



Sensitivity to ε_S and ε_T

FUTURE

$\varepsilon_{S,T}$ @ $\mu = 2 \text{ GeV (MS-bar)}$



$$d\Gamma \sim \Gamma_0 (1 + b m_e / E_e)$$

Prospective beta decay measurements competitive with strong LHC constraints, probing $\Lambda_{S,T} \sim 10 \text{ TeV}$

Note: $\lesssim 10\%$ uncertainty on $g_{S,T}$ is essential to achieve competitiveness

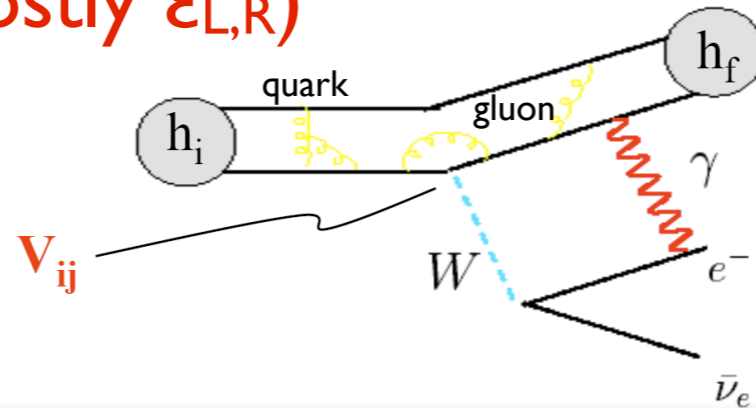
Current low-E data:
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 $b(0^+ \rightarrow 0^+)$, $A(n)$

**Gonzalez-Alonso,
Naviliat-Cuncic,
Severijns, 1803.08732**

**Bhattacharya et al
1806.09006**

How do we probe the ε_α ? (2)

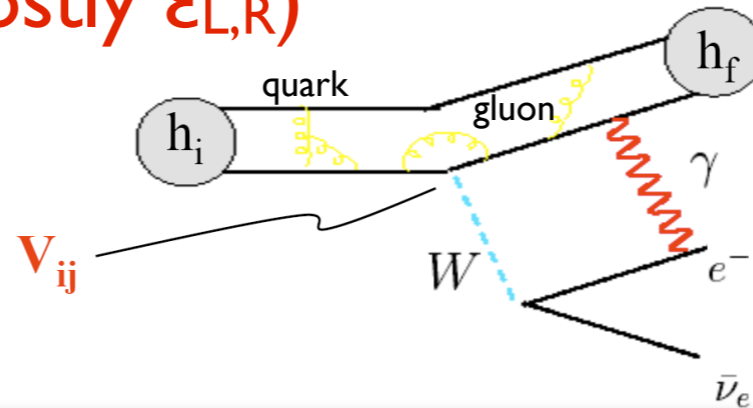
2. Total decay rates (mostly $\varepsilon_{L,R}$)



$$\Gamma_k = (G_F^{(\mu)})^2 \times |\bar{V}_{ij}|^2 \times |M_{\text{had}}|^2 \times (1 + \delta_{RC}) \times F_{\text{kin}}$$

How do we probe the ε_α ? (2)

2. Total decay rates (mostly $\varepsilon_{L,R}$)



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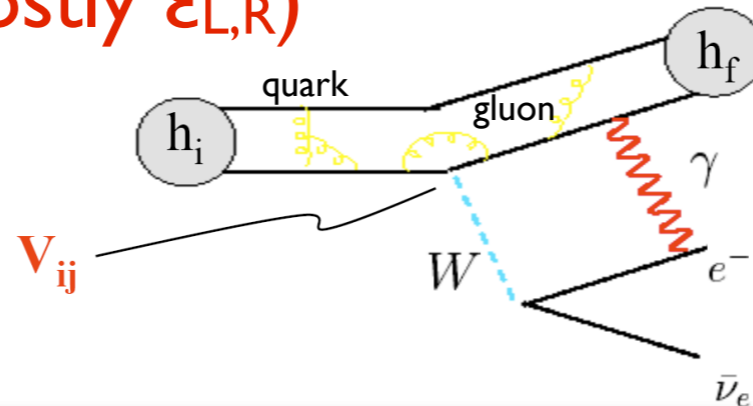
Lifetimes,
BRs

Experimental input

Q-values \rightarrow
phase space

How do we probe the ε_α ? (2)

2. Total decay rates (mostly $\varepsilon_{L,R}$)



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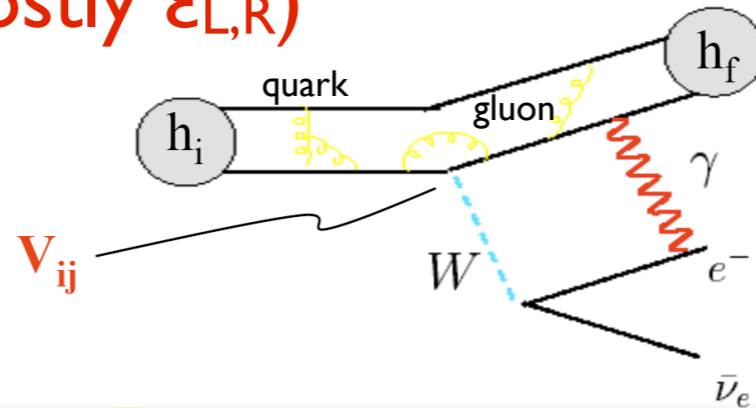
Theory input

Hadronic / nuclear
matrix elements
and radiative corrections

LQCD, chiral EFT,
dispersion relations
+ expt. input (g_A)

How do we probe the ϵ_α ? (2)

2. Total decay rates (mostly $\epsilon_{L,R}$)



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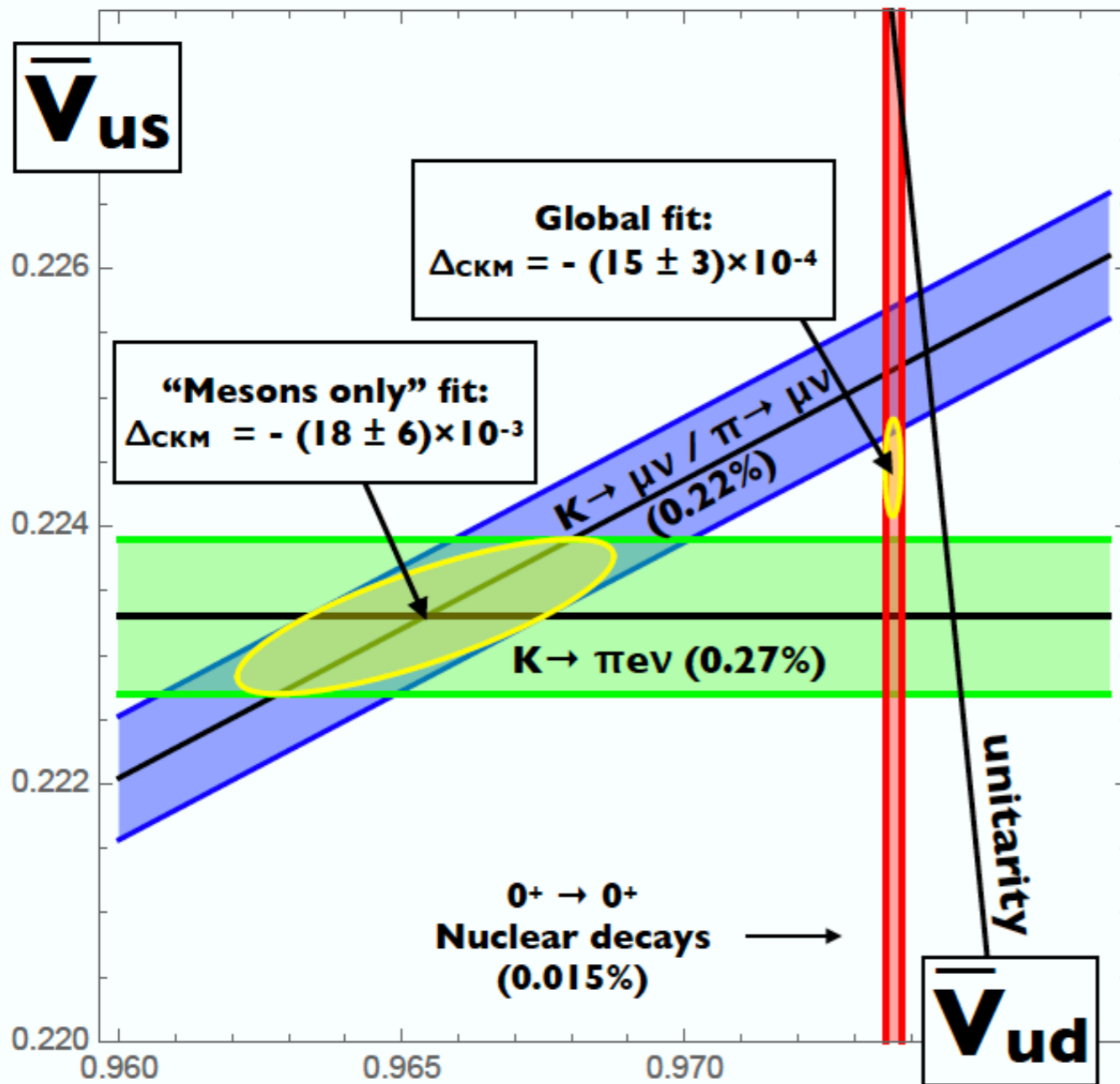
Channel-dependent
effective CKM element

$$\bar{V}_{ud} = V_{ud} \left[1 + \epsilon_L \pm \epsilon_R + b(\epsilon_S, \epsilon_T) \tilde{F}_{\text{kin}} \right]$$

$$|\bar{V}_{ud}|^2 + |\bar{V}_{us}|^2 + |\bar{V}_{ub}|^2 = 1 + \Delta_{\text{CKM}}(\epsilon_i)$$

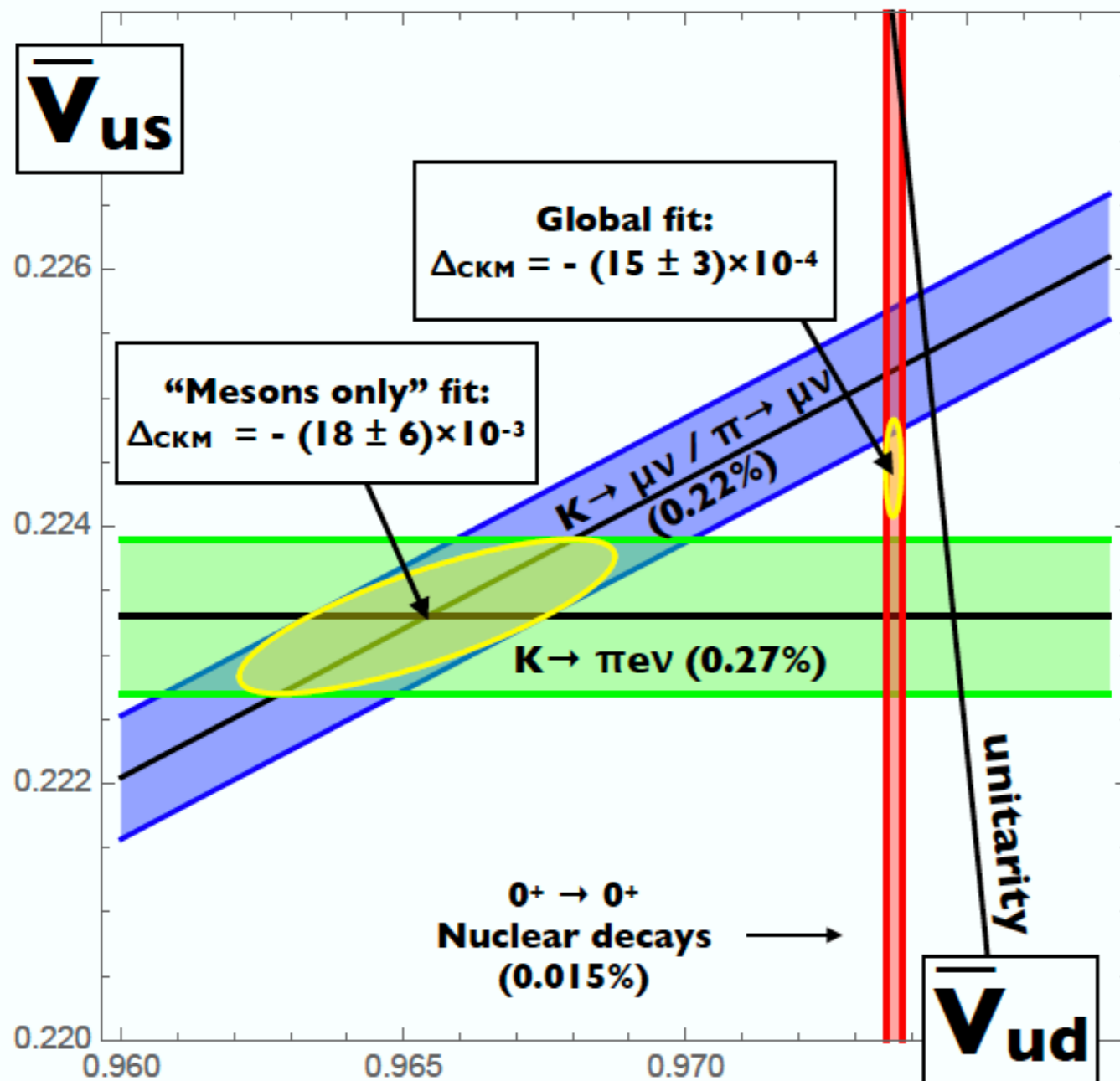
CKM unitarity test

$$|\bar{V}_{ud}|^2 + |\bar{V}_{us}|^2 + |\cancel{\bar{V}_{ub}}|^2 = 1 + \Delta_{\text{CKM}}(\epsilon_i)$$



CKM unitarity test

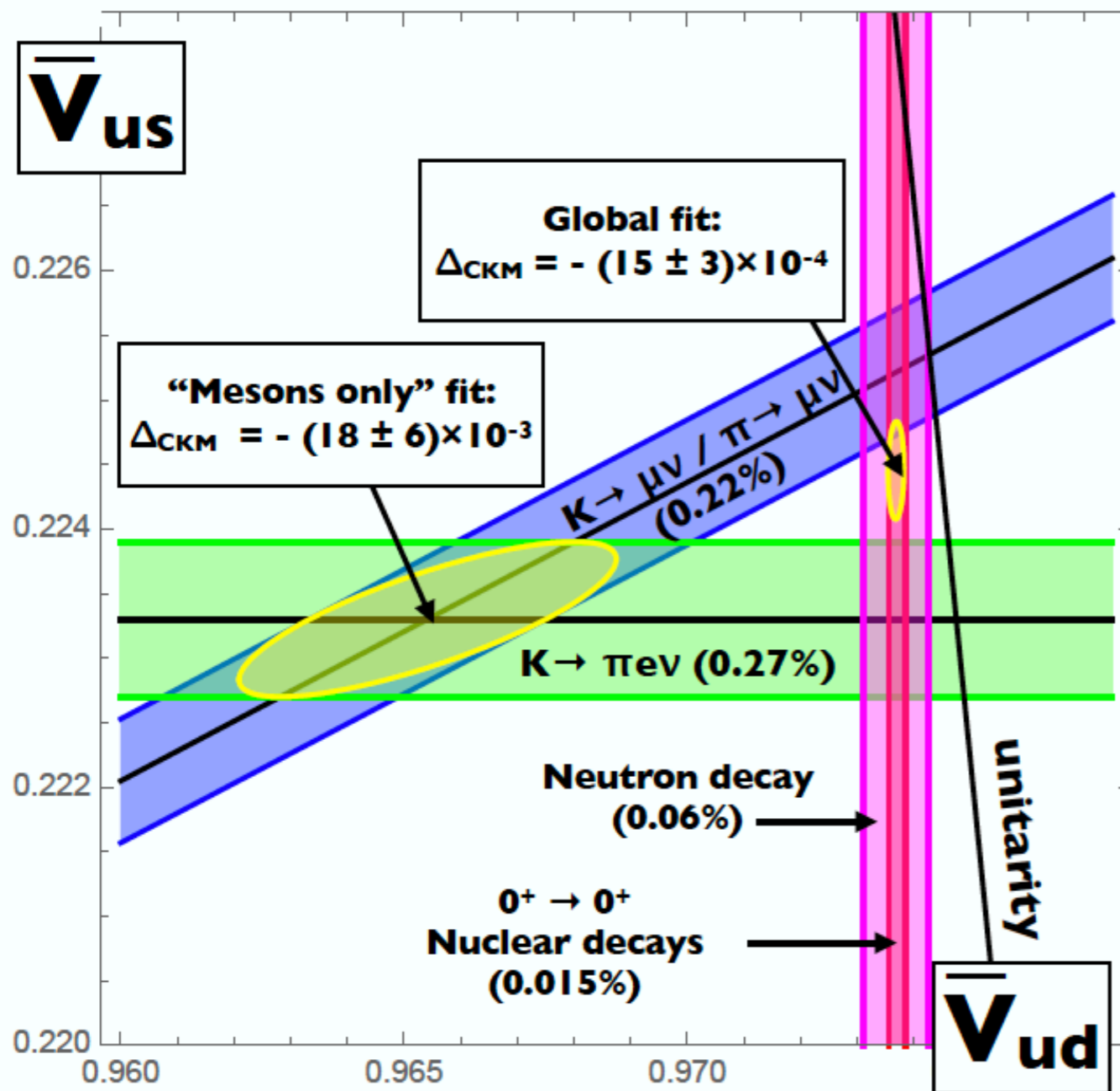
$$|\bar{V}_{ud}|^2 + |\bar{V}_{us}|^2 + |\bar{V}_{ub}|^2 = 1 + \Delta_{\text{CKM}}(\epsilon_i)$$



- Discrepancy could be explained in terms of $\epsilon_R^{(s)}$ ($\sim 0.4\%$), and ϵ_R ($\sim 0.1\%$)
- Used radiative corrections from Seng et al, 1807.10197. Discrepancy goes from $5\sigma \rightarrow 3\sigma$ if use Czarnecki et al, 1907.06737 → **Importance of model-independent treatments of radiative corrections, in all decay channels (requires EFT + lattice QCD).**
- Theory analysis of nuclear decays at 0.015% level currently suffers from nuclear structure uncertainties: **neutron will be the arbiter**

Impact of neutron measurements

$$|\bar{V}_{ud}|^2 + |\bar{V}_{us}|^2 + \cancel{|\bar{V}_{ub}|^2} = 1 + \Delta_{\text{CKM}}(\epsilon_i)$$



- V_{ud} from neutron decay currently at 0.06%
- Independent extraction of V_{ud} @ 0.015%, via neutron decay requires:

$$\bar{V}_{ud} = \left[\frac{4908.6(1.9) \text{ s}}{\tau_n (1 + 3\bar{g}_A^2)} \right]^{1/2}$$

$$\delta\tau_n \sim 0.7\text{s} \rightarrow 0.3\text{s} \rightarrow 0.1\text{s}$$

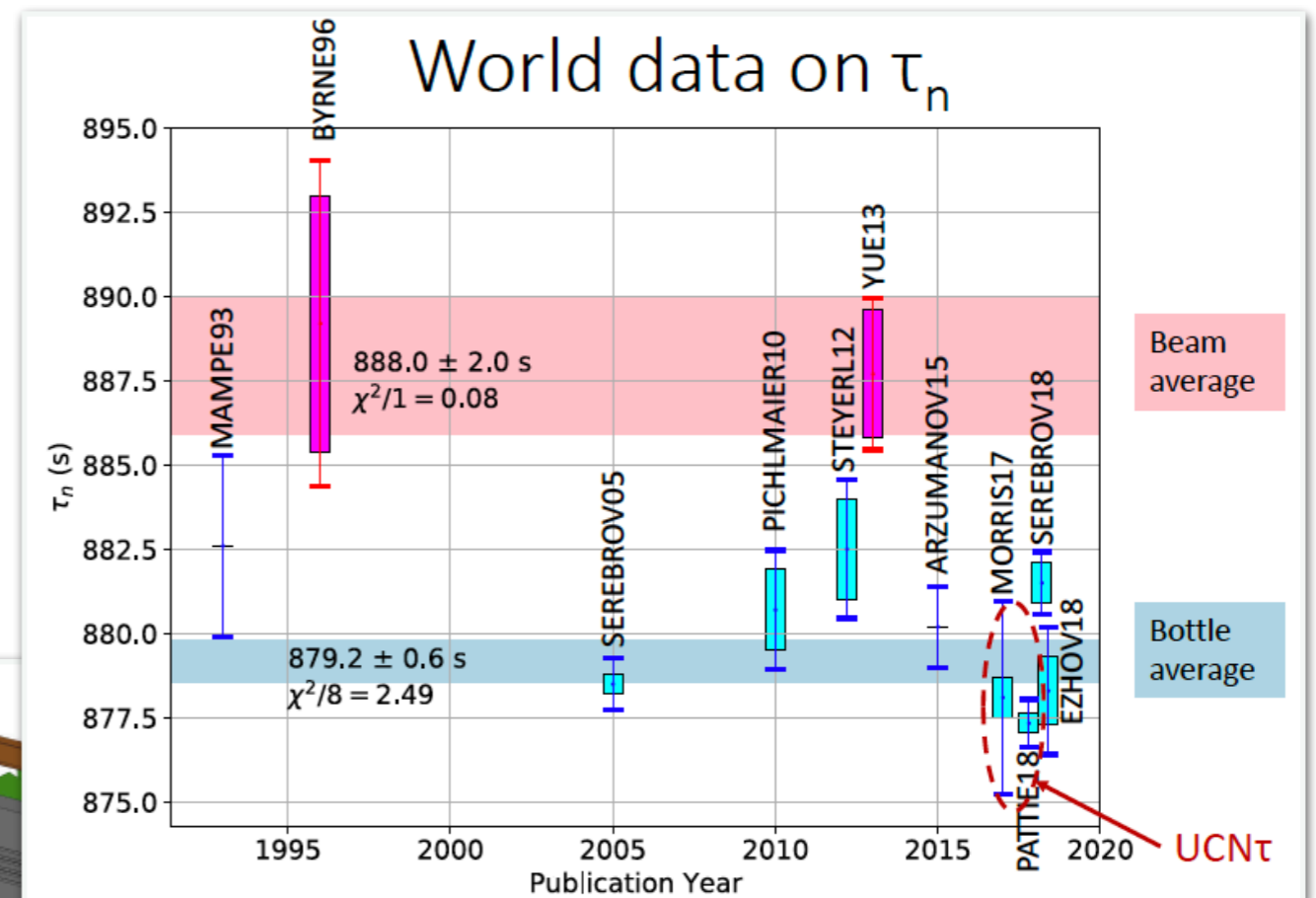
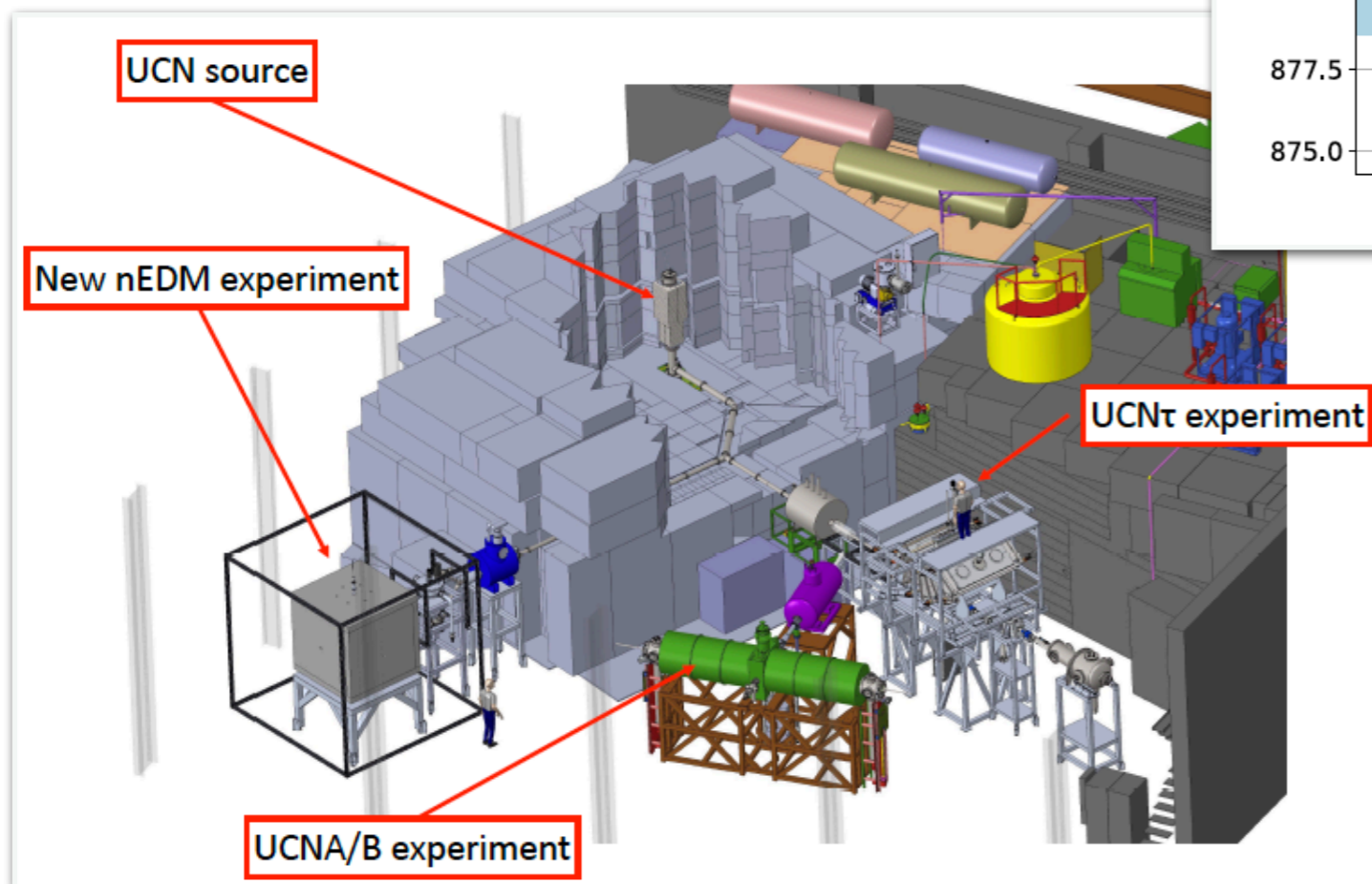
@ LANL UCNT \rightarrow UCNT+

$$\delta g_A/g_A \sim 0.044\% \rightarrow 0.02\%$$

@ LANL UCNA+ \rightarrow PERC

Impact of LANL

- UCN τ result is the world's most precise (Science, 2018)

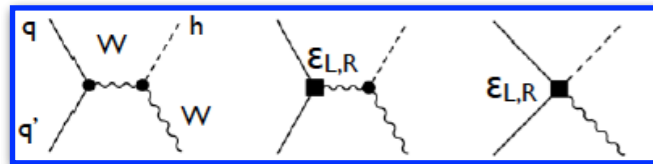


- New LDRD investment:
 - $\delta\tau_n \rightarrow 0.1$ s
 - $\delta g_A/g_A \rightarrow 0.03\%$
 - Rad. Corr.: LQCD + EFT

Sensitivity to ϵ_L and ϵ_R

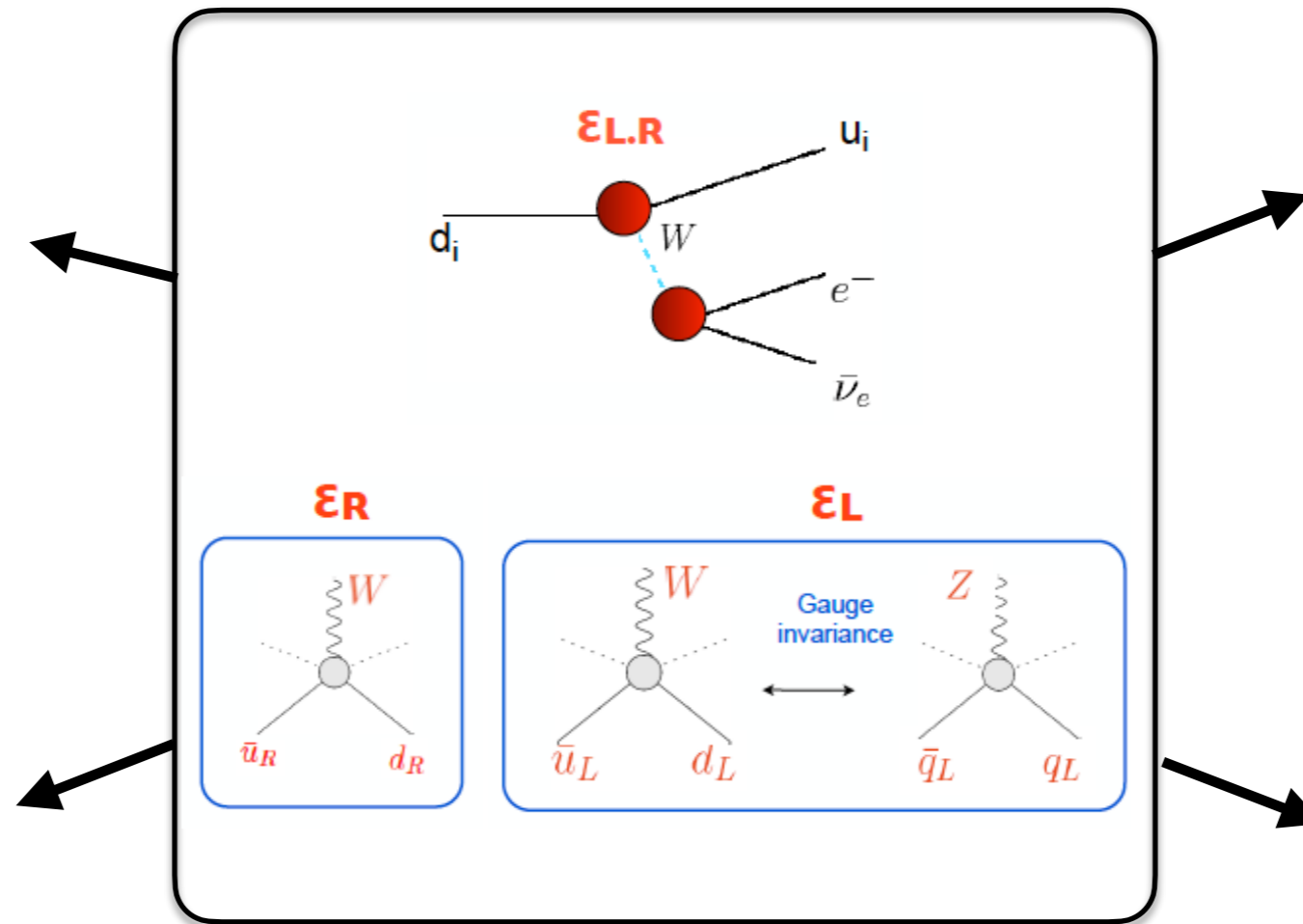
S. Alioli, VC, W. Dekens, J. de Vries, E. Mereghetti 1703.04751

Associated Higgs
production at LHC



LEP, SLC

Z-pole:
Falkowski et al
1706.03783



Neutron decay:
 $\lambda = g_A (1 - 2 \epsilon_R)$

Constraint on ϵ_R uses
 $g_A = 1.271(13)$
(CalLat 1805.12030)

$\Delta_{\text{CKM}} \propto \epsilon_L + \epsilon_R$

$\delta\Gamma_{(\pi \rightarrow \mu \nu)} \propto \epsilon_L - \epsilon_R$
[f_π from LQCD]

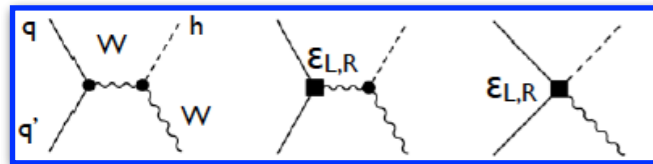
Due to weak
isospin symmetry, vertex corrections
involve the Higgs & Z bosons

Sensitivity to ϵ_L and ϵ_R

S. Alioli, VC, W. Dekens, J. de Vries, E. Mereghetti 1703.04751

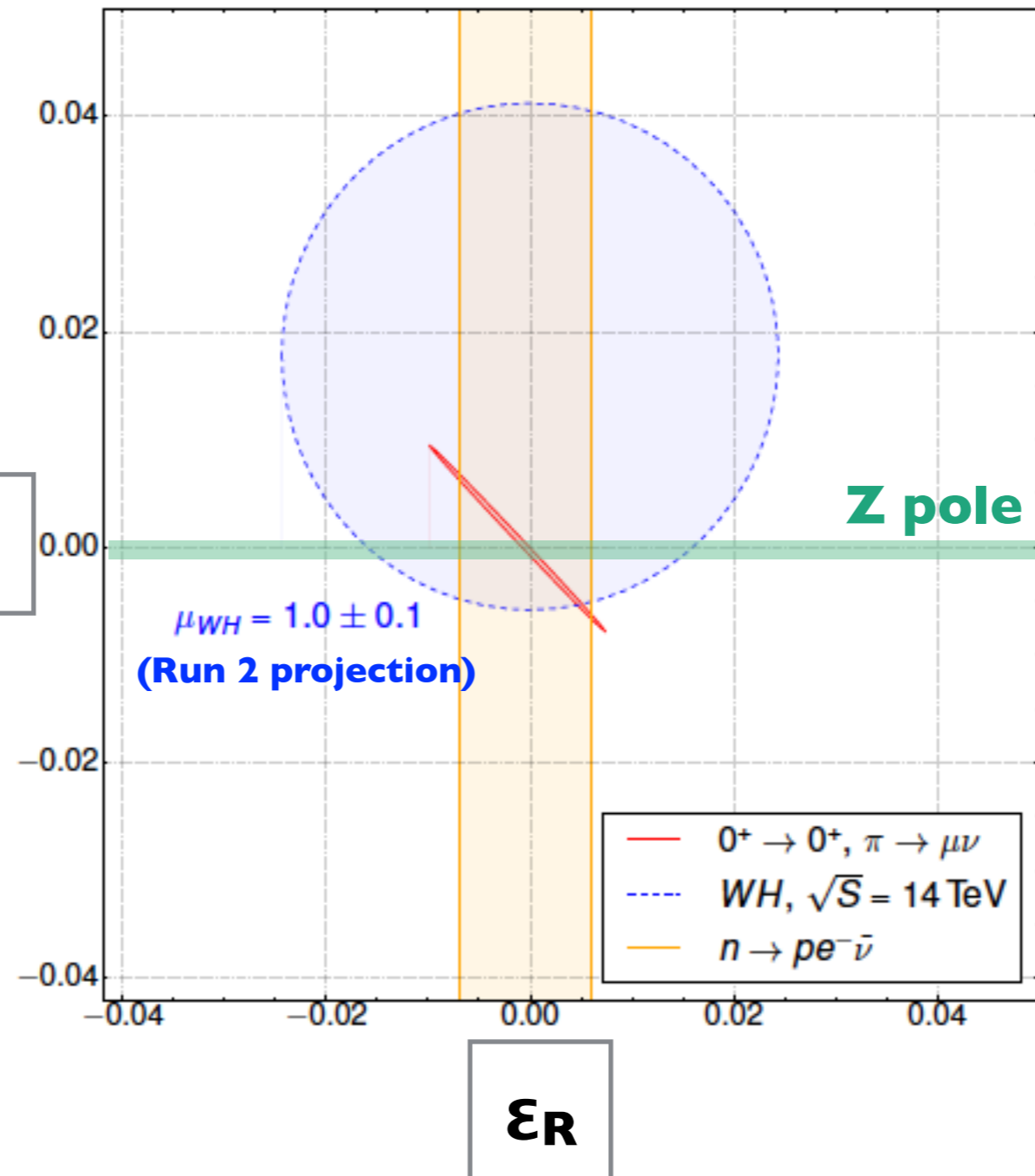
90%CL allowed regions, assumes only two operators at high scale

Associated Higgs
production at LHC



ϵ_L

LEP, SLC
Z-pole:
Falkowski et al
1706.03783



Neutron decay:
 $\lambda = g_A (1 - 2 \epsilon_R)$

Constraint on ϵ_R uses
 $g_A = 1.271(13)$
(CalLat 1805.12030)

$\Delta_{\text{CKM}} \propto \epsilon_L + \epsilon_R$

$\delta\Gamma_{(\pi \rightarrow \mu\nu)} \propto \epsilon_L - \epsilon_R$
[f_π from LQCD]

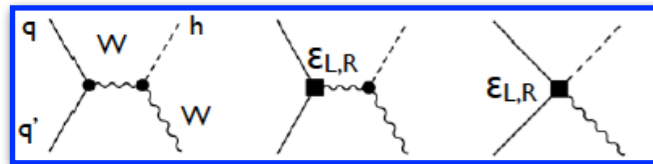
ϵ_R

Sensitivity to ϵ_L and ϵ_R

S. Alioli, VC, W. Dekens, J. de Vries, E. Mereghetti 1703.04751

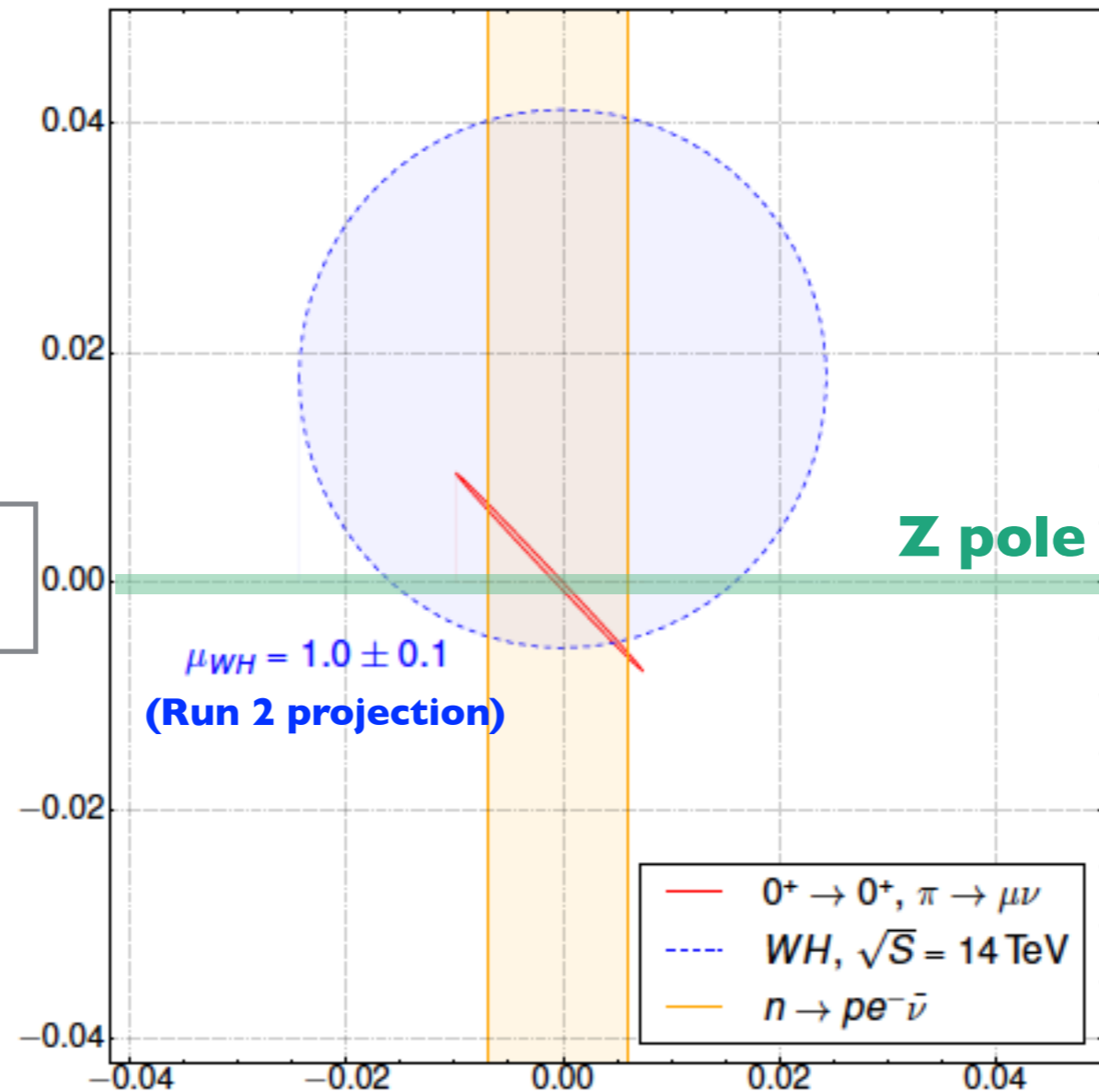
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$\Delta_{CKM} \propto \epsilon_L + \epsilon_R$

$\delta\Gamma_{(\pi \rightarrow \mu \nu)} \propto \epsilon_L - \epsilon_R$
[f_π from LQCD]

β decays more constraining than collider: probing $\Lambda_{L,R} \sim 20$ TeV

β decays summary

- EFT shows that a discovery window exists well into the LHC era
 - Beta decays play unique role in probing vertex corrections ϵ_L - ϵ_R (unmatched sensitivity compared to LHC)
 - Beta decays can be competitive in probing scalar and tensor interactions if precision reaches $< 0.1\%$ (ϵ_S - ϵ_T plots)
- LANL key player in the international neutron physics scene

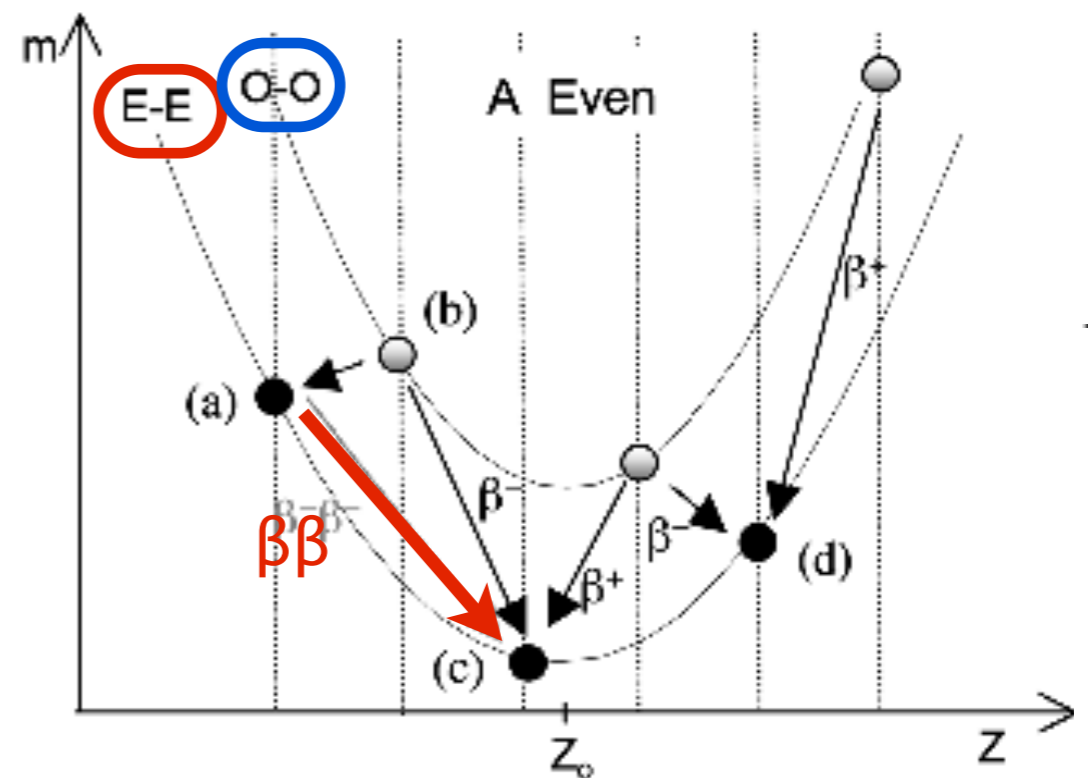
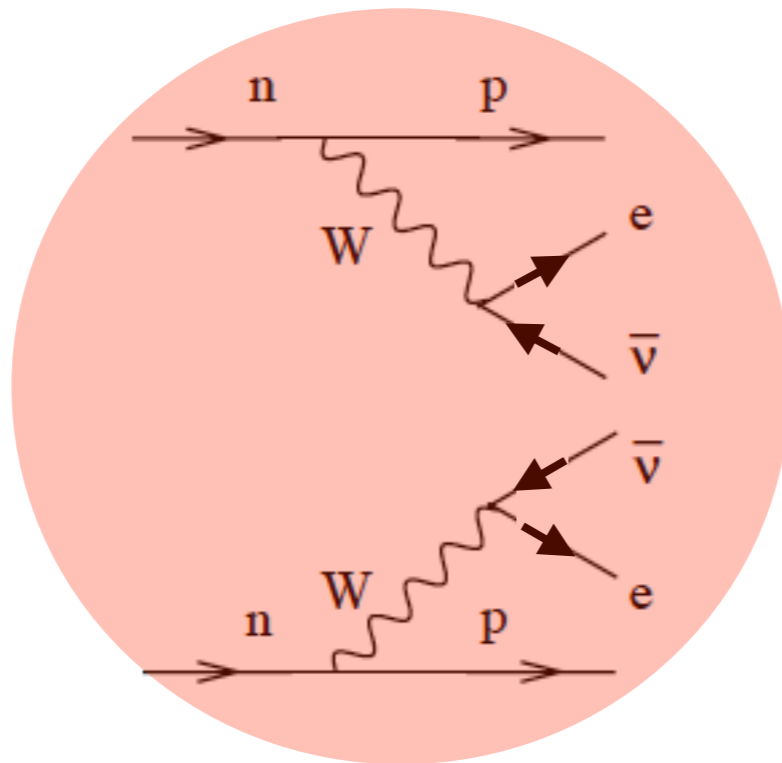
Neutrinoless double beta decay

Double beta decay

- For certain even-even nuclei (^{48}Ca , ^{76}Ge , ^{136}Xe , ...), single β decay is energetically forbidden $\rightarrow \beta\beta$ decay!



M. Goppert Mayer, 1935



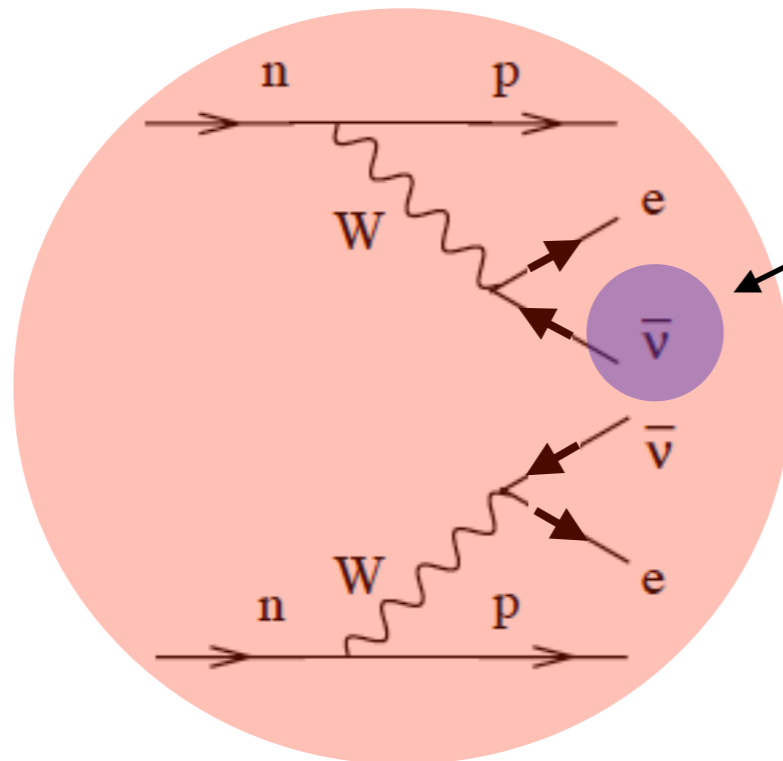
- $2\nu\beta\beta$ is the rarest process ever observed, with $T_{1/2} \sim 10^{21}$ years (first observation in 1987)

Neutrinoless double beta decay?

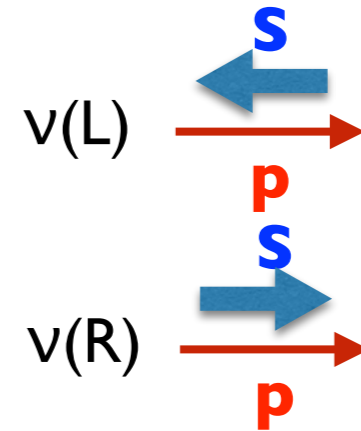
- Yes, if neutrinos are Majorana particles (i.e. their own antiparticles)



W. H. Furry, 1939



This is just $\nu(R)$, which mixes with $\nu(L)$ via mass insertion



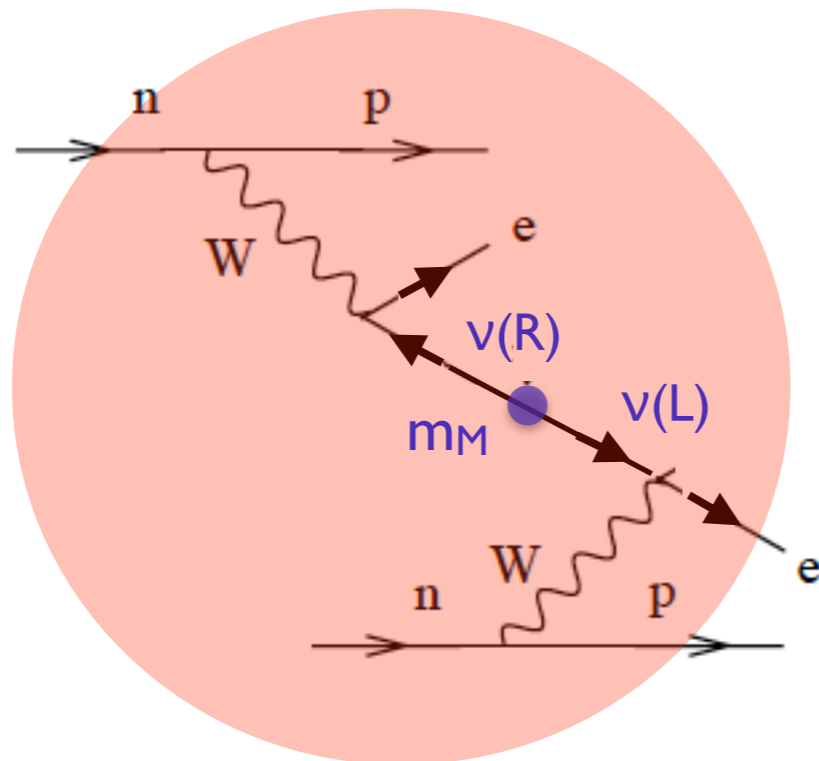
E. Majorana, 1937

Neutrinoless double beta decay?

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W. H. Furry, 1939



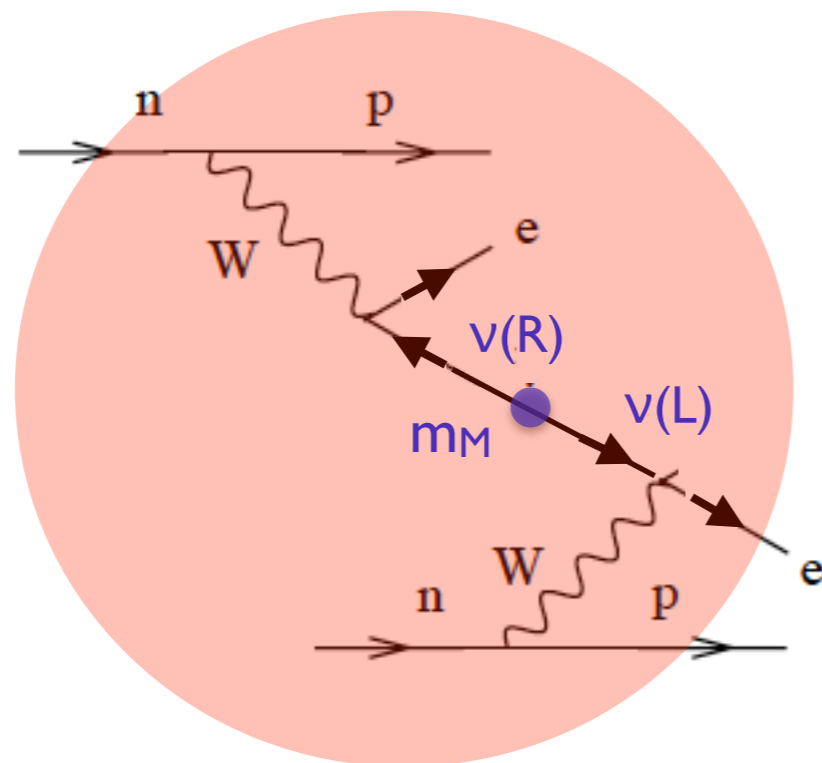
“Subject to the usual limitations on the meaning of such language, one can say that a (virtual) neutrino is emitted together with one of the electrons and reabsorbed when the other electron is emitted.”

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“Subject to the usual limitations on the meaning of such language, one can say that a (virtual) neutrino is emitted together with one of the electrons and reabsorbed when the other electron is emitted.”

- Key point: in $0\nu\beta\beta$ Lepton Number changes by two units.
 ν_M exchange is just one possible mechanism. Furry understood this:

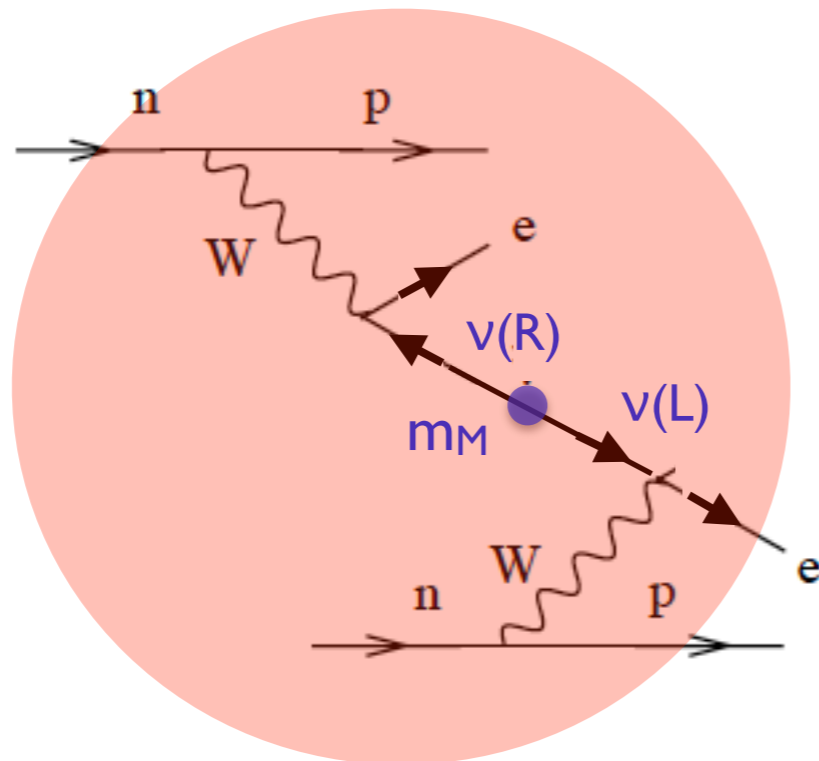
“The Majorana form of the theory is not the only one that permits this new form of disintegration [...]. The Majorana theory provides, so to speak, a canonical form.”

Neutrinoless double beta decay?

- Yes, if neutrinos are Majorana particles (i.e. their own antiparticles)

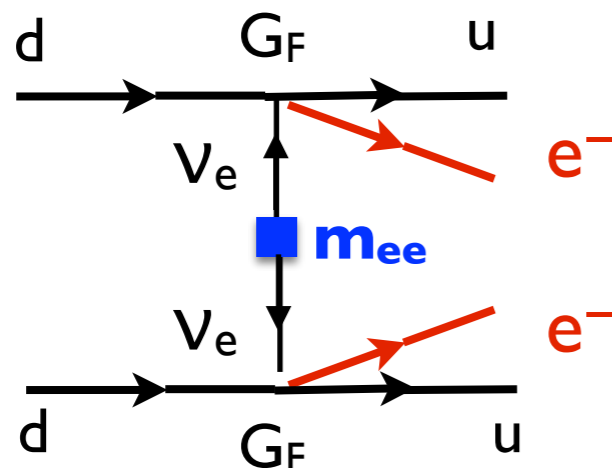


W. H. Furry, 1939

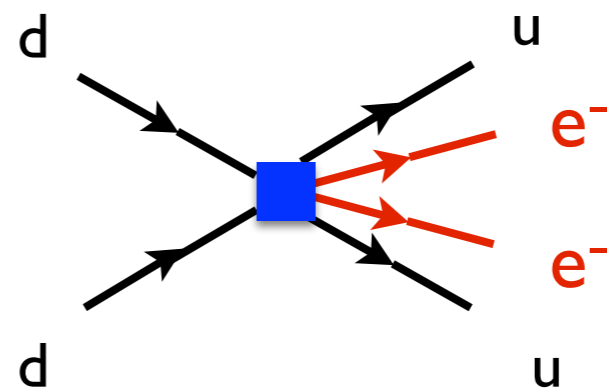


“Subject to the usual limitations on the meaning of such language, one can say that a (virtual) neutrino is emitted together with one of the electrons and reabsorbed when the other electron is emitted.”

- Modern viewpoint on LNV:



but also

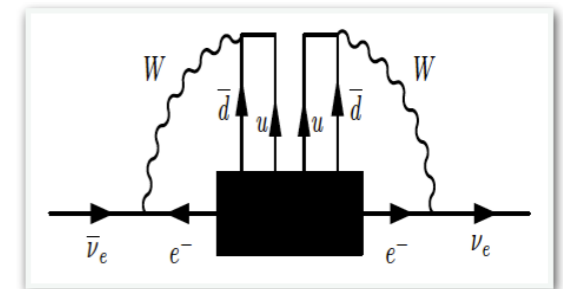


Exchange of heavier neutrinos or other Majorana particles. At low-energy induce six-fermion operator $\sim 1/\Lambda^5$

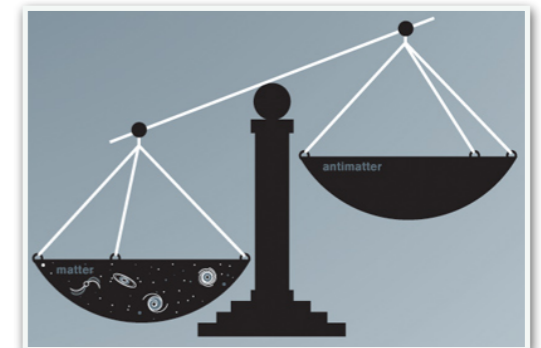
Significance of $0\nu\beta\beta$

- B-L conserved in SM $\rightarrow 0\nu\beta\beta$ = new physics, with far-reaching implications

- Demonstrate that neutrinos are their own antiparticles
- Probe origin of neutrino mass
- Establish L non-conservation, a key ingredient to generate the baryon asymmetry via leptogenesis



Shechter-
Valle 1982



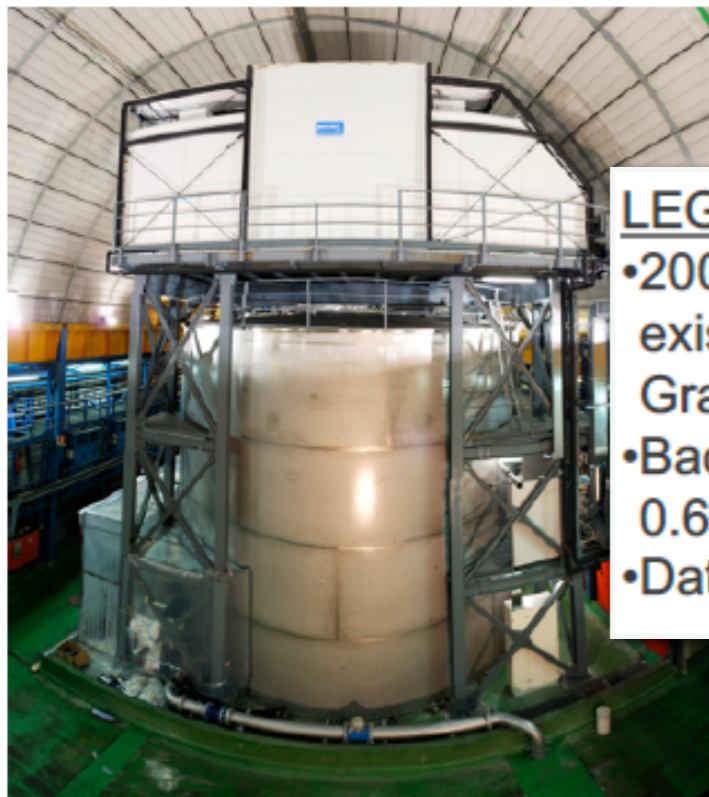
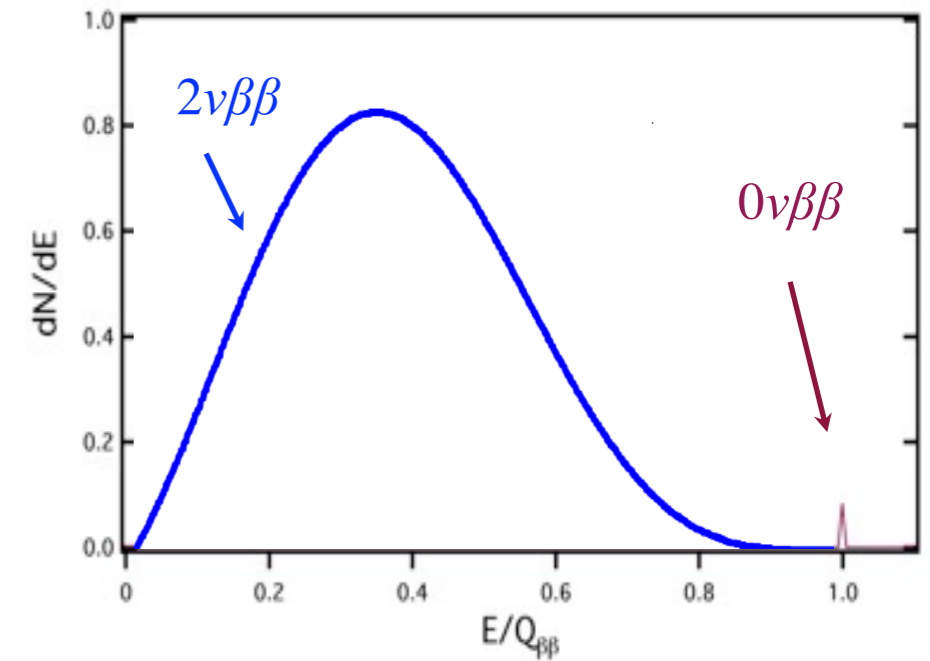
Fukujita-
Yanagida 1987

Construction of ton-scale $0\nu\beta\beta$ experiment ($T_{1/2} > 10^{27-28}$ yr) is the top priority for new project starts in the 2015 NSAC Long Range Plan

The quest is on...

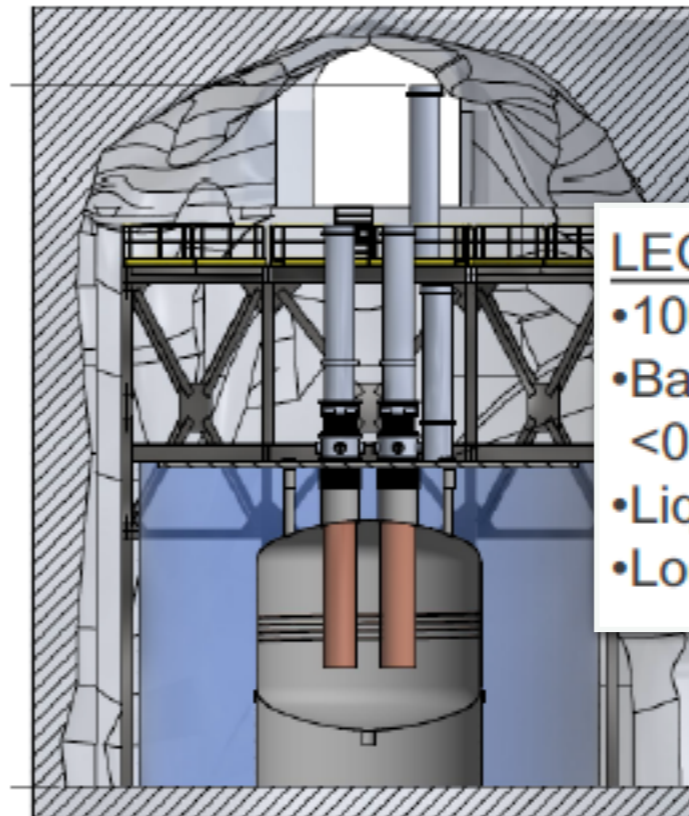
- Several experiments worldwide
- LANL co-leads ^{76}Ge -based search.
Majorana Demonstrator \rightarrow LEGEND

~ 50 institutions, ~ 250 scientists, S. Elliott co-spokesperson



LEGEND-200:

- 200 kg in upgrade of existing infrastructure at Gran Sasso
- Background goal $0.6 \text{ cts}/(\text{FWHM t yr})$
- Data start ~ 2021



LEGEND-1000:

- 1000 kg, staged via 4 payloads
- Background goal $< 0.03 \text{ cts}/(\text{FWHM t yr})$
- Liquid Ar cryostat shield
- Location to be selected

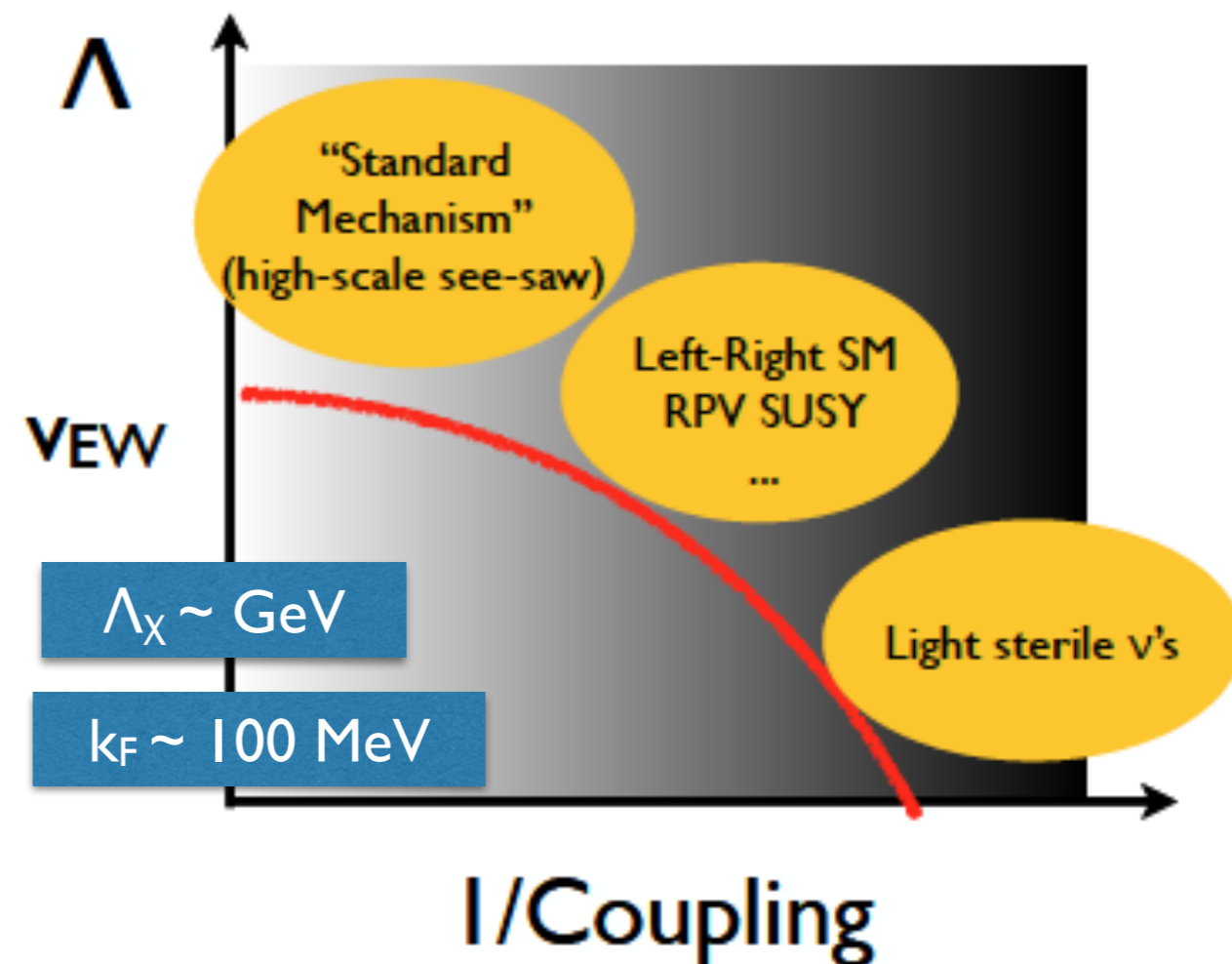
$T_{1/2} \sim 10^{28} \text{ yr}$

The quest is on...

- Challenging experiments and challenging theory!

- Connect *any* source of LNV to nuclei: multi-scale problem!

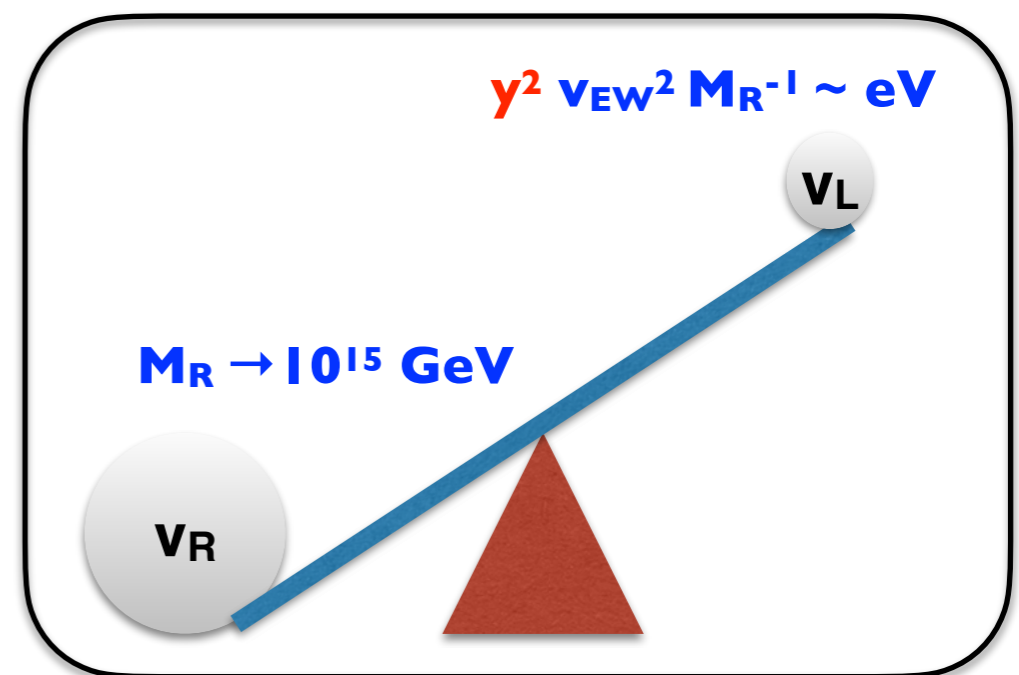
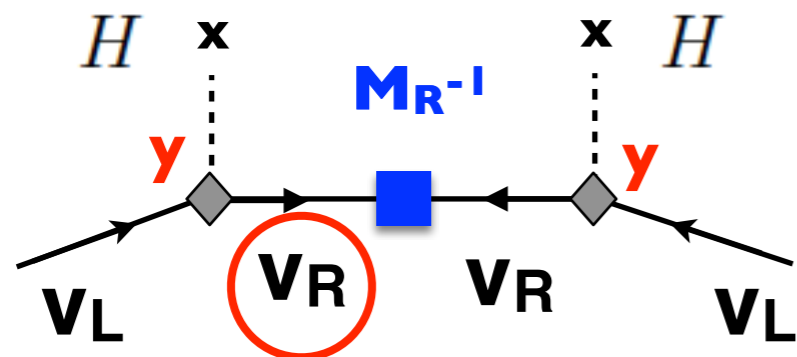
- Our approach: use effective field theory techniques to bridge scales, combined with LQCD and nuclear structure



$$T_{1/2} \sim (m_W/\Lambda)^A (\Lambda_\chi/m_W)^B (k_F/\Lambda_\chi)^C$$

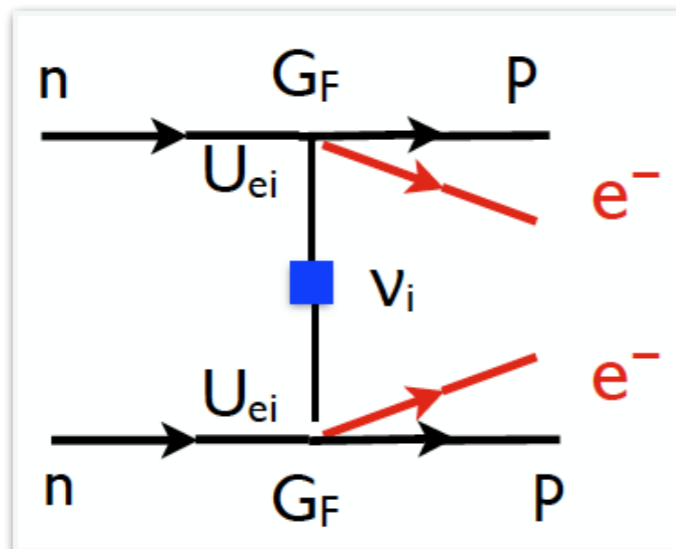
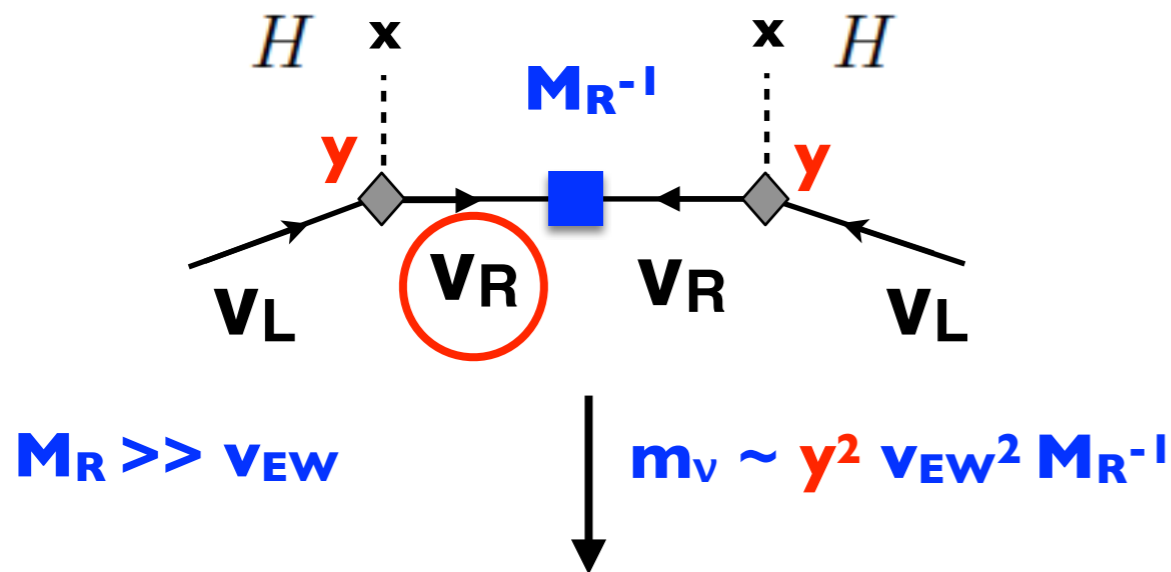
High-scale seesaw

- Majorana mass generated by exchange of heavy neutrinos, neutral under all SM charges (=sterile)



High-scale seesaw

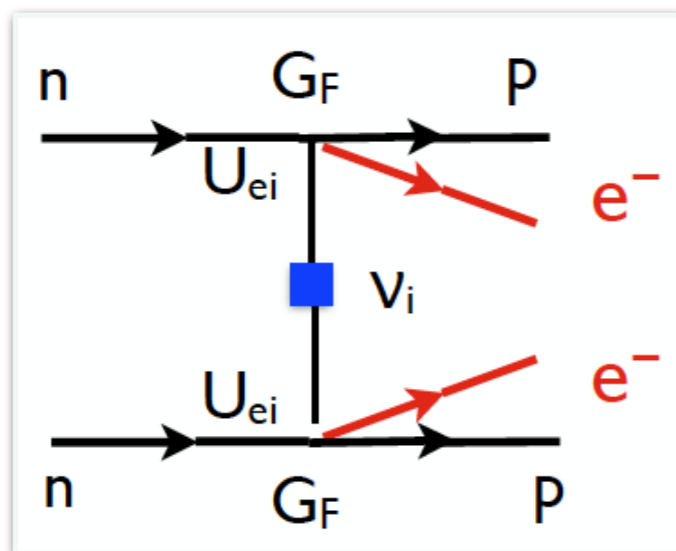
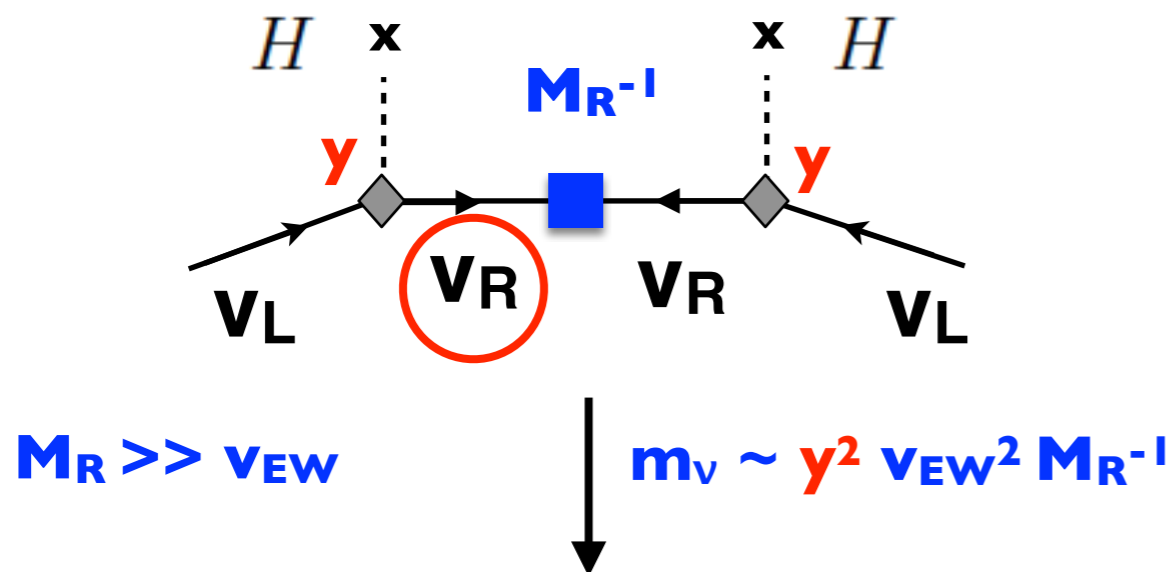
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$0\nu\beta\beta$ mediated by light neutrinos

High-scale seesaw

- Majorana mass generated by exchange of heavy neutrinos, neutral under all SM charges (=sterile)



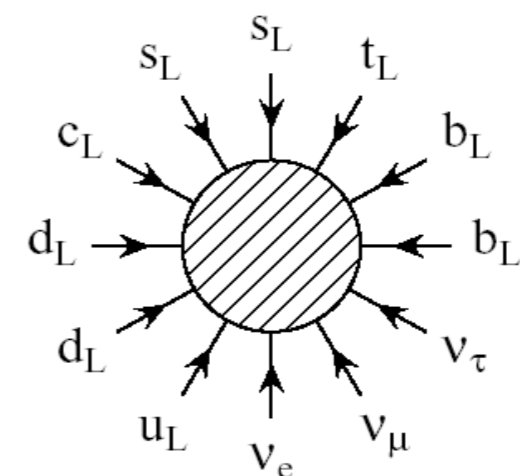
$0\nu\beta\beta$ mediated by light neutrinos

Baryogenesis via Leptogenesis

1) CP- and L- violating
out-of-equilibrium
decays of heavy $\nu_{Ri} \Rightarrow n_L$

$$\Gamma(\nu_R \rightarrow H^* \ell) \neq \Gamma(\nu_R \rightarrow H \bar{\ell})$$

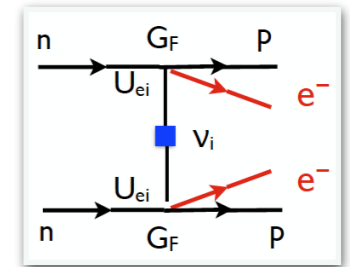
2) EW sphalerons $\Rightarrow n_B = \# n_L$



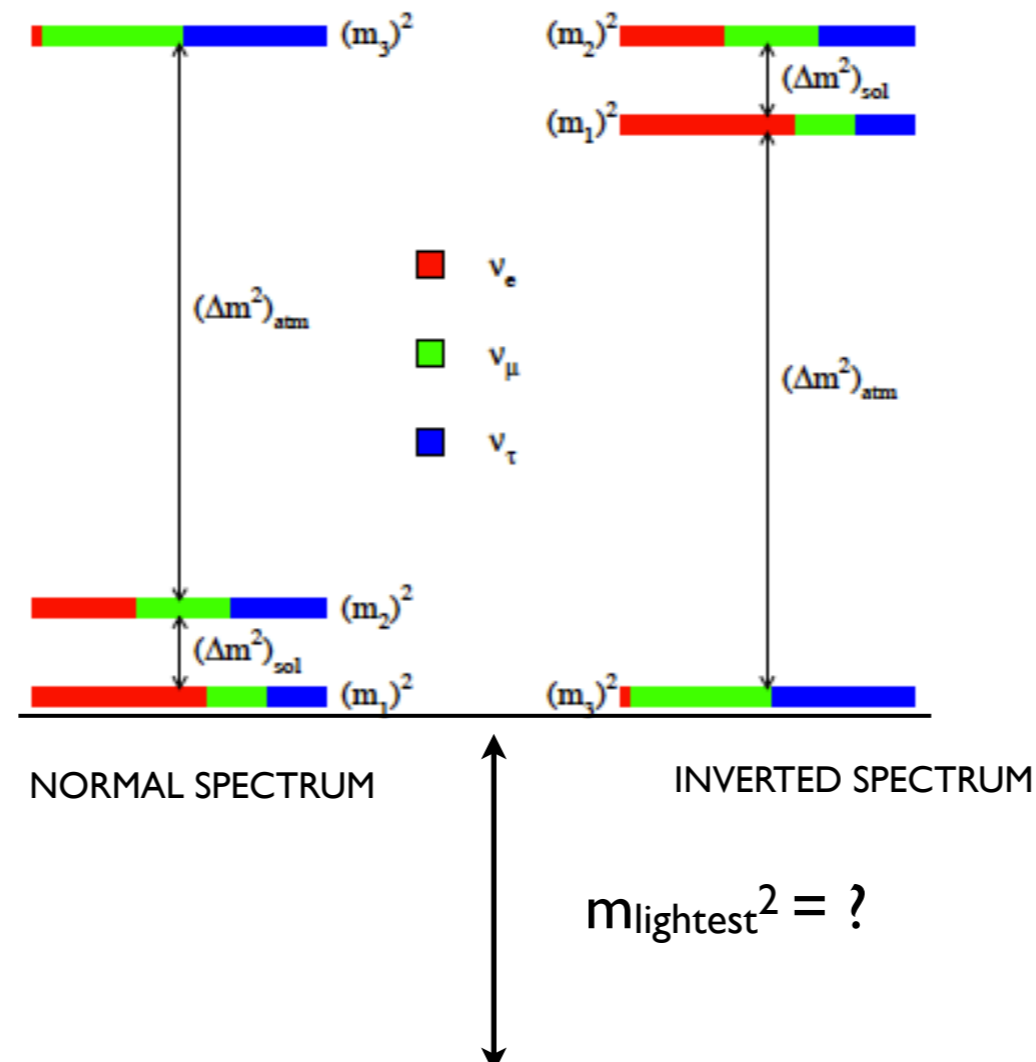
High-scale seesaw: discovery potential

- In this case $0\nu\beta\beta$ is a *direct* probe of ν Majorana mass: $\Gamma \propto |M_{0\nu}|^2 (m_{\beta\beta})^2$

$$\langle m_{\beta\beta} \rangle^2 = \left| \sum_i U_{ei}^2 m_{\nu i} \right|^2$$



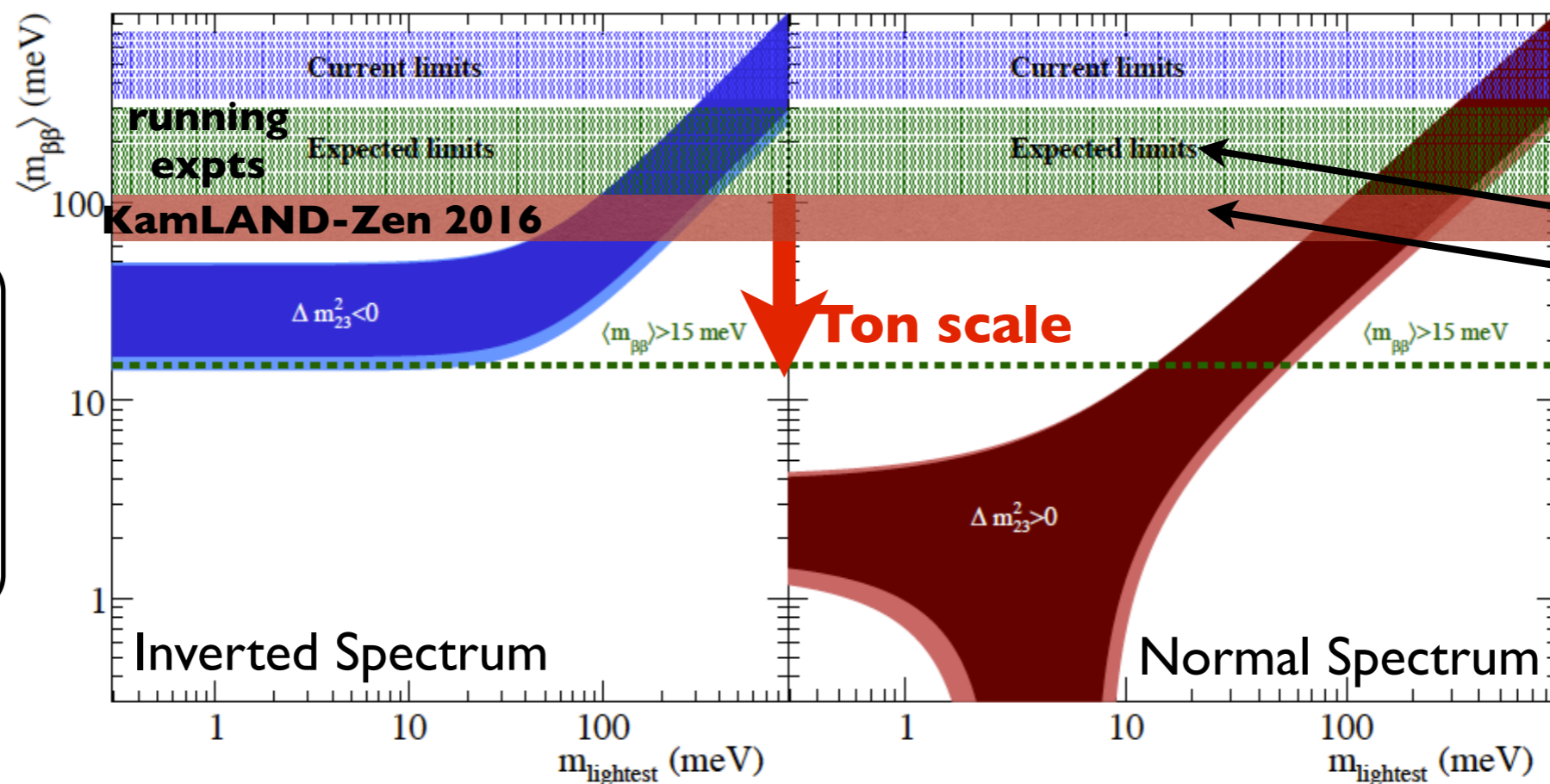
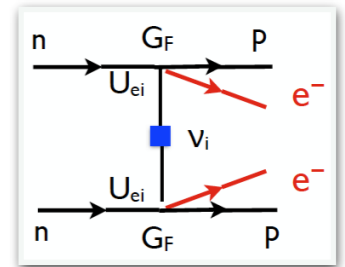
Strong correlation with
oscillation parameters



High-scale seesaw: discovery potential

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Dark bands:
unknown phases

Light bands:
uncertainty from
oscillation
parameters(90% CL)

Assume range for
nuclear matrix
elements from
different many-body
methods

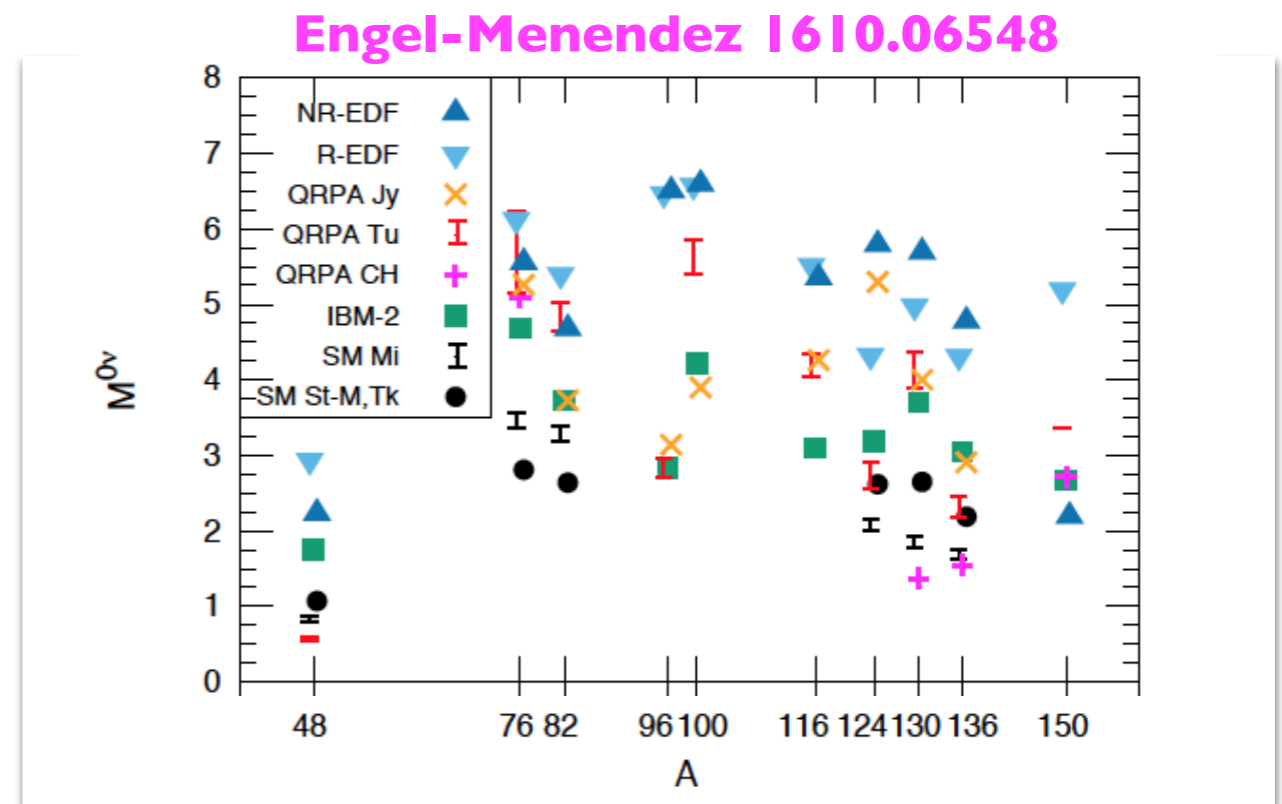
Plot by K. Heeger

Discovery @ ton-scale possible for **inverted spectrum** or **$m_{\text{lightest}} > 50 \text{ meV}$**

High-scale seesaw: theory developments

- Sensitivity to $m_{\beta\beta}$ affected by large uncertainty in “nuclear matrix elements”:

$$\Gamma \propto |M_{0\nu}|^2 (m_{\beta\beta})^2$$

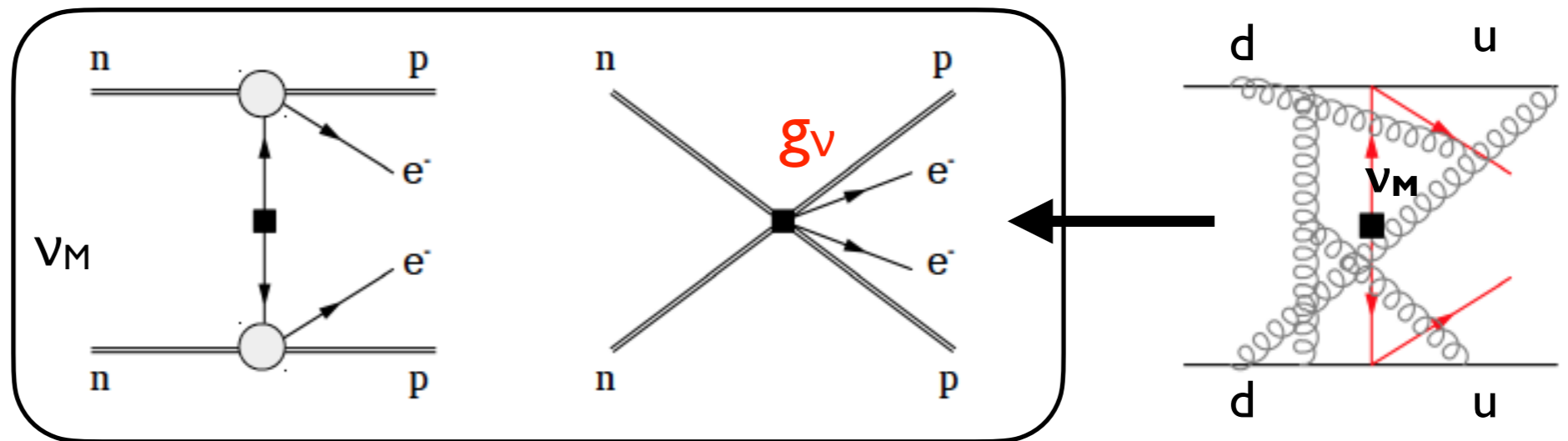


- Steps towards controllable uncertainties in matrix elements:
 - Use chiral EFT as guiding principle
 - Lattice QCD for *hadronic* matrix elements (e.g. $nn \rightarrow pp$)
 - “Ab initio” nuclear structure calculations:
from light nuclei (benchmark) to ^{48}Ca and ^{76}Ge

New insights from EFT

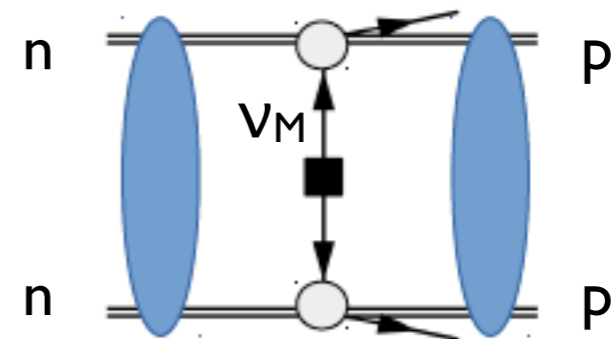
V. Cirigliano, W. Dekens, J. de Vries, M. Graesser, E. Mereghetti, S. Pastore, U. van Kolck
1802.10097, Phys.Rev.Lett. 120 (2018) no.20, 202001

- Transition operator to leading order in Q/Λ_χ ($Q \sim k_F \sim m_\pi$, $\Lambda_\chi \sim \text{GeV}$)



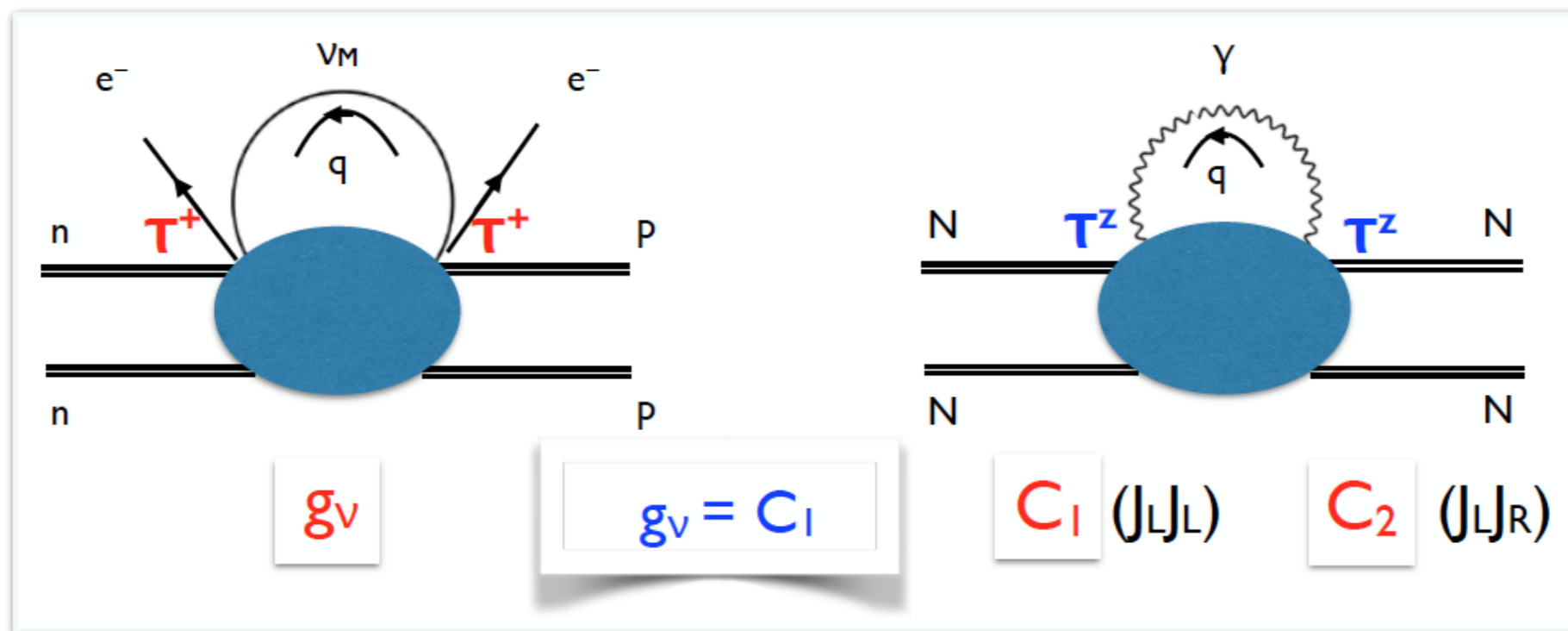
‘Usual’ V_M exchange $\sim 1/Q^2$
Coulomb-like potential

‘New’: short-range coupling $g_v \sim 1/Q^2$.
Required by renormalization of $nn \rightarrow ppee$ amplitude
in presence of strong interactions



Connection with data?

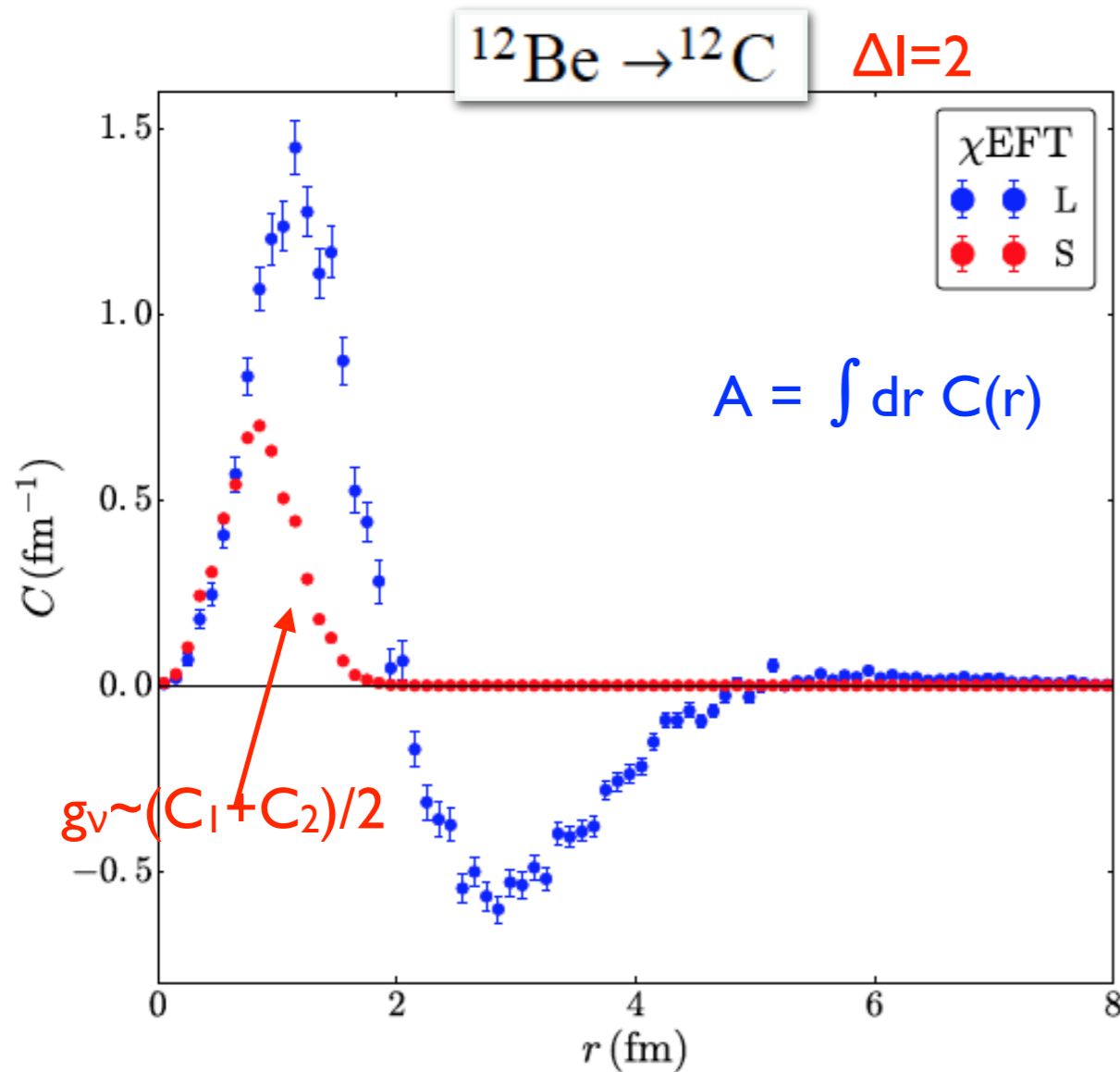
- Chiral symmetry relates g_V to one of two $I=2$ EM LECs (hard γ 's & ν 's)



- NN scattering data determine $(C_1 + C_2)$, but not $g_V = C_1$
- Assuming $g_V \sim (C_1 + C_2)/2$, what is the impact on $m_{\beta\beta}$ extraction?

Impact on nuclear matrix elements

V.C. , W. Dekens, J. de Vries, M. Graesser, E. Mereghetti, S. Pastore, M. Piarulli, U. van Kolck, R. Wiringa, 1907.11254

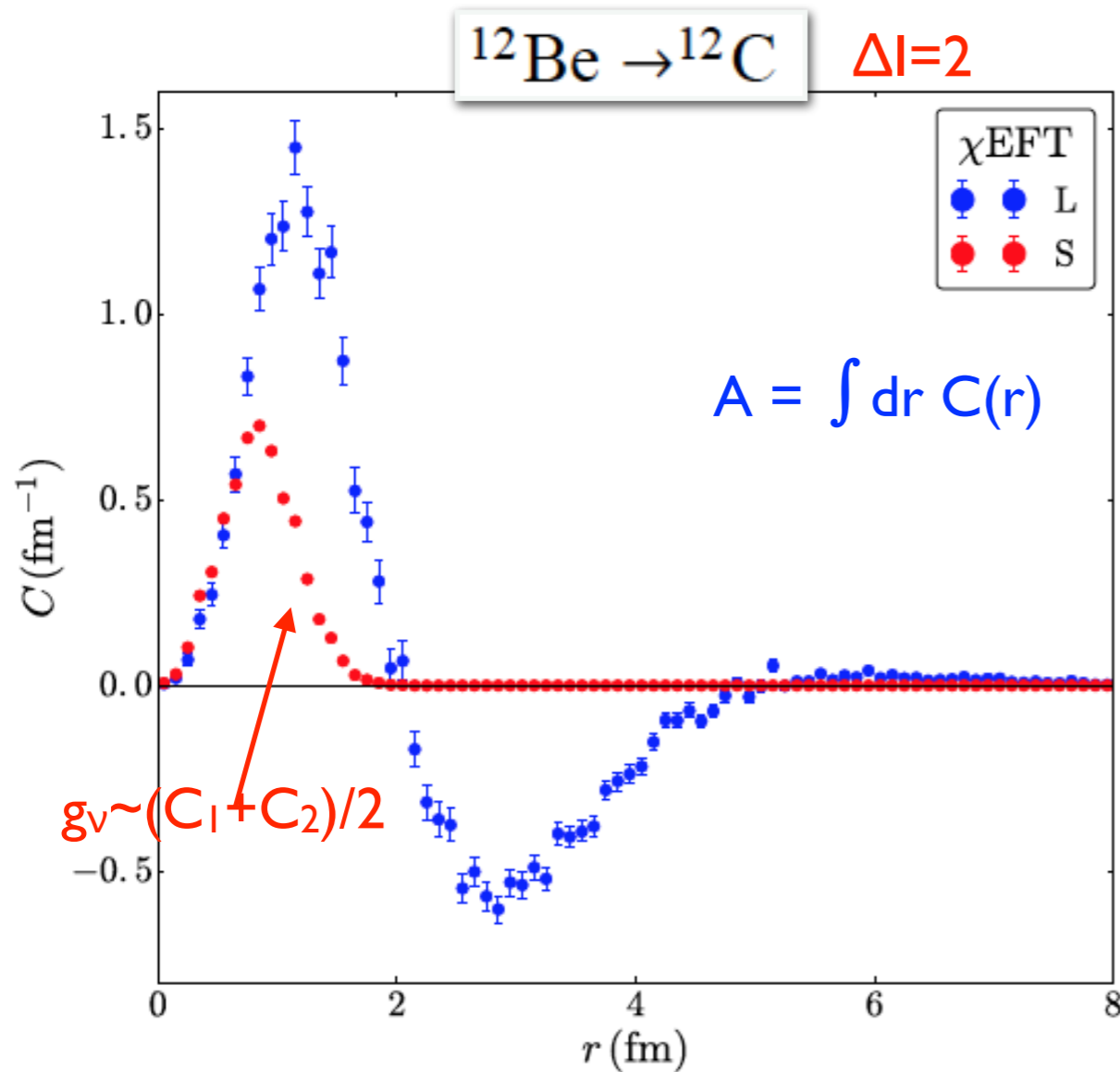


Evaluate **impact in light nuclei** using VMC wavefunctions from Norfolk chiral potential [1606.06335]

g_v contribution sizable in $\Delta I=2$ transition (due to node):
for $A=12$, $A_S/A_L = 0.75$

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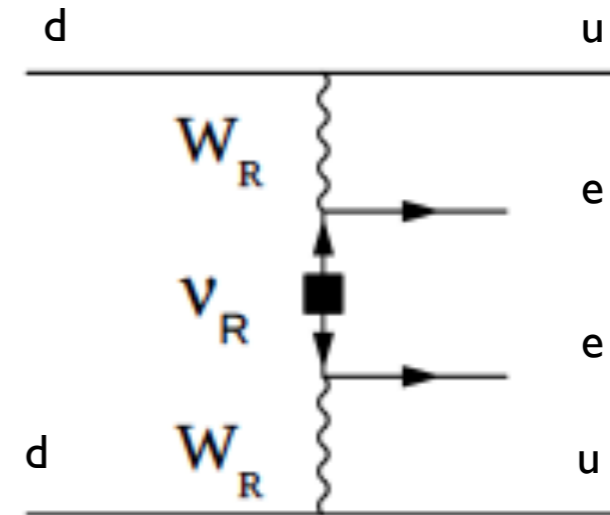
g_v contribution sizable in $\Delta I=2$ transition (due to node):
for $A=12$, $A_S/A_L = 0.75$

Transitions of experimental interest ($^{76}\text{Ge} \rightarrow ^{76}\text{Se}, \dots$) have $\Delta I=2$ (and node) \Rightarrow expect significant effect!

Determination of g_v is a 'decadal' challenge: analytic methods & lattice QCD

TeV-scale LNV

- **TeV sources of LNV** may lead to observable contributions to $0\nu\beta\beta$ *not directly related to the exchange of light neutrinos*



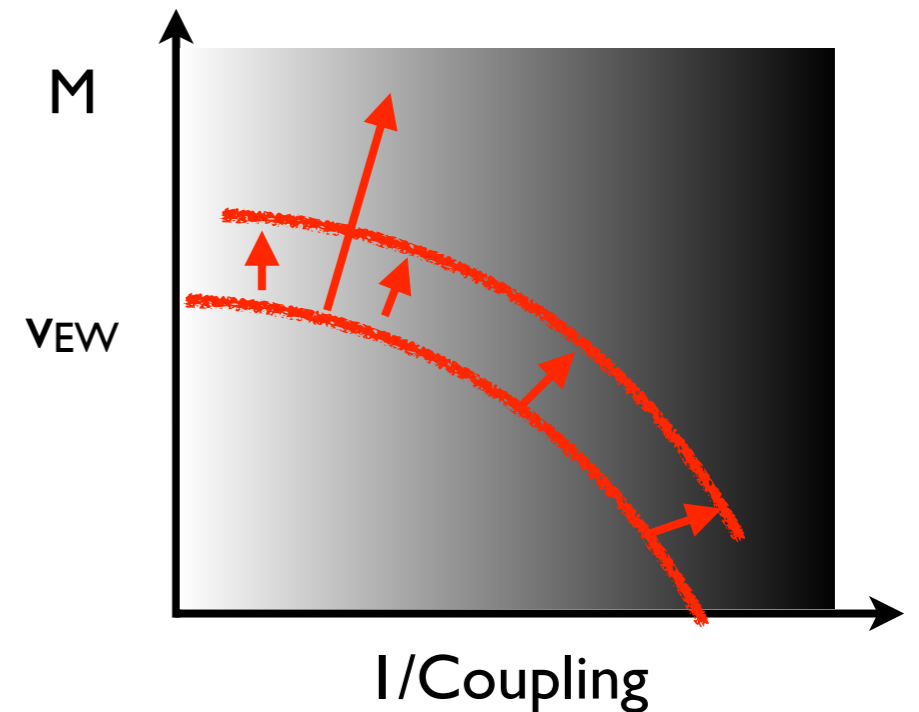
- May lead to correlated (or precursor!) signal at LHC: $pp \rightarrow ee jj$
- New contributions *can interfere with* $m\beta\beta$ or add incoherently, significantly affecting the interpretation of experimental results

$0\nu\beta\beta$ summary

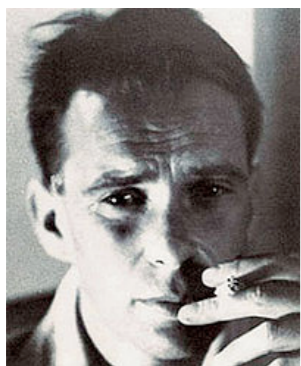
- Ton-scale $0\nu\beta\beta$ searches ($T_{1/2} > 10^{27-28}$ yr) have **great discovery potential** — we simply don't know the origin of neutrino mass and the scale Λ associated with LNV
- Exciting prospects to improve theory uncertainties thanks to synergy of **EFT**, **lattice QCD**, and **nuclear structure**
- LANL co-leads a world wide ^{76}Ge -based experimental search

Concluding comments

- Precision frontier experiments are exploring uncharted territory in our search for new physics. **Important component of DOE mission (HEP, NP)**
- LANL at the forefront of the precision frontier through **experiment, theory, high performance computing**
- Illustrated challenges and impact through β and $\beta\beta$ decays



Thank you!



A drawing by
Bruno Tuschek